

Study of the Dependence of the End-Point Bremsstrahlung Energy on the Residual Gas Pressure during the Pyroelectric Source Operation in Vacuum

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Abstract—Experimental results on the determination of the end-point bremsstrahlung energy during the operation of a pyroelectric source in vacuum at different residual gas pressures are presented. The range of operating residual gas pressures is determined.

Keywords: pyroelectric deflector, X-ray radiation

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The operation scheme of the pyroelectric X-ray source based on two crystals was proposed for the first time by Danon [1]. In the first test experiments, Danon detected X-ray generation with a record at that time end-point energy of ~160 keV. The X-ray penetrability directly depends on its end-point energy. Previously [2], the dependence of the maximum X-ray energy generated during pyroelectric source operation in vacuum on the preliminary change in the pyroelectric source temperature was studied. In [3], the dependence of the X-ray yield on the temperature variation rate of a single crystal source is demonstrated. The effect of the target shape in the pyroelectric source based on a single crystal on X-ray spectral characteristics was studied as well [4]. This paper is devoted to the determination of the optimum pressure at which the end-point energy of X-ray bremsstrahlung generated during the two-crystal pyroelectric source operation reaches a maximum. It should be noted that the two-crystal X-ray source being the object of the present study, in contrast to the single crystal source [2], can generate X-rays with doubled end-point energy. Figure 1 shows the schematic diagram of the experimental setup for studying the dependence of the end-point energy of X-ray bremsstrahlung on the residual gas pressure in vacuum during the pyroelectric source operation.

The study was carried out in vacuum chamber 5 45 L in volume. The residual gas pressure in the vacuum chamber is measured using an ERSTVAK MTM9D-KF25 vacuum gauge; the pressure is changed by a vacuum shut off valve of the MDC E-GV-4000M manual drive. Two pyroelectric lithium niobate LiNbO₃ crystals 1 10 × 20 × 20 mm in size are placed in the vacuum chamber 5 at a distance of 11 mm coaxially and in parallel to each other. The spontaneous polarization vectors of each pyroelectric are codirected. Crystals 1 are heated by silicon semiconductor diodes MUR 1660 3 through heat conductors 2 made of aluminum 40 mm in diameter and 5 mm thick. The operating supply current of semiconductor silicon diodes 3 was 4.5 A. The temperature of pyroelectric crystals 1 was measured by a K-type thermocouple 4 mounted on the rear surface of heat conductors 2. Each assembly consisting of a pyroelectric crystal 1, heat conductor 2, semiconductor silicon diode 3, and K-type thermocouple 4 was fixed on opposite walls of the vacuum chamber 5. Pyroelectric crystals 1 were cooled naturally due to heat dissipation on heat conductors 2 and vacuum chamber walls 5. The X-ray spectrum was measured using an Amptek CdTe 123 semiconductor X-ray detector 6 with a working area of 25 mm² and a thickness

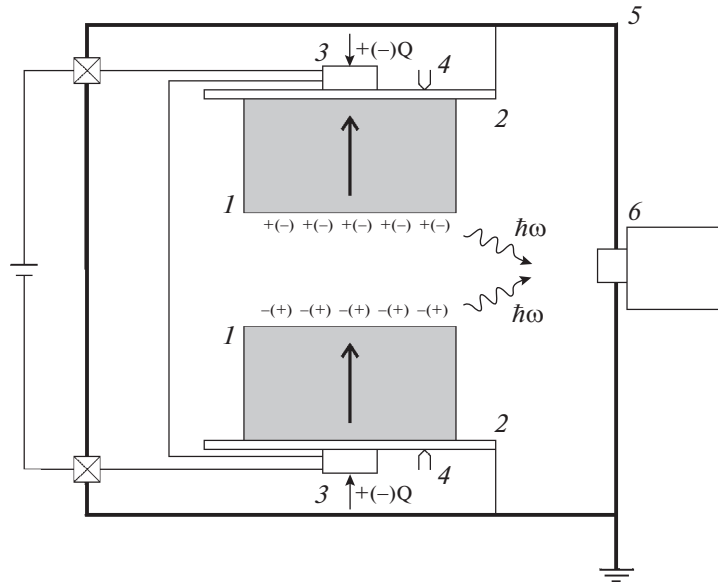


Fig. 1. Schematic of the experimental setup: (1) pyroelectric lithium niobate crystals, (2) heat conductor, (3) silicon semiconductor diodes, (4) K-type thermocouple, (5) vacuum chamber, and (6) X-ray detector.

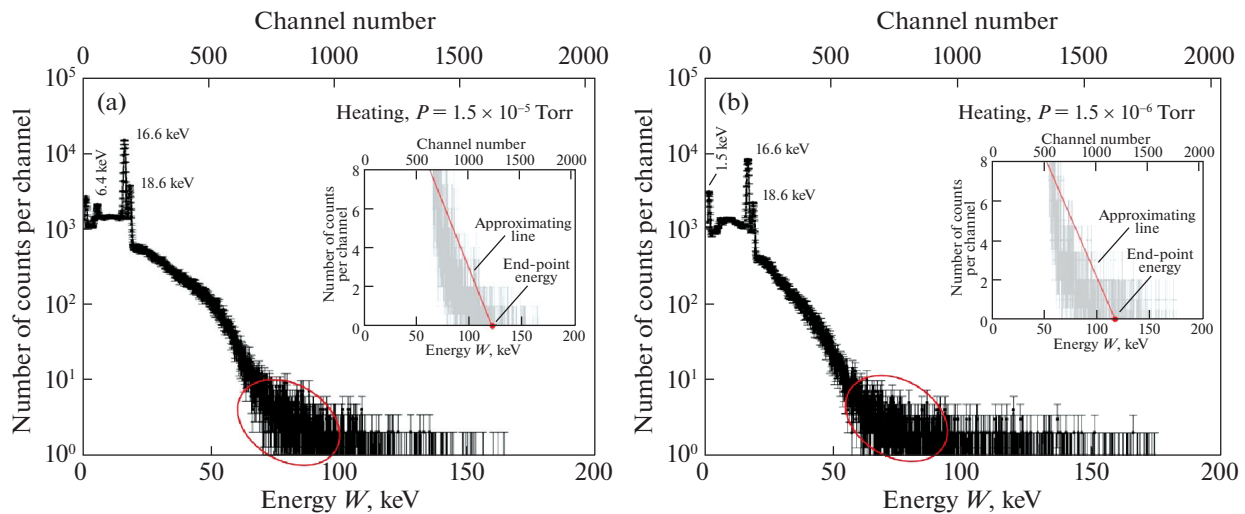


Fig. 2. X-ray spectra generated during (a) heating and (b) cooling of the pyroelectric source at a residual gas pressure of 1.5×10^{-5} Torr.

of 1 mm. The entrance window of the semiconductor detector is made of beryllium foil 100 μm thick. The detector was calibrated by spectral lines of the ^{57}Co cobalt isotope. The detector energy resolution was 530 eV in the ^{57}Co spectral peak with an energy of 14.4 keV; its peaking time was 1 μs .

A system of vacuum pumps provides a pressure of 1.3×10^{-6} Torr in the vacuum chamber 5. Pyroelectric crystals 1 are heated from room temperature of 25°C to 71.2°C within 900 seconds. At a supply current of 4.5 A in each semiconductor diode 3, the average temperature variation rate of crystals 1 is $3^\circ\text{C}/\text{min}$. Then, pyroelectric crystals 1 are naturally cooled due to thermal dissipation on heat conductors 2 and vacuum chamber walls 5 until the temperature of each crystal reached 30°C . The crystal natural cooling in vacuum to room temperature requires significant time-consuming (several hours); therefore, the temperature variation range of pyroelectric crystals was changed by $30\text{--}71.2^\circ\text{C}$.

The X-ray spectra were measured during heating and cooling of the pyroelectric deflector at various residual gas pressures from 1.3×10^{-6} to 1.5×10^{-2} Torr with a step depending on the manual adjustment accuracy of the vacuum shut off valve. Figure 2a shows the X-ray spectrum measured during heating of

Table 1. Results of processing of obtained experimental data

P , Torr	W , keV (heating)	W , keV (cooling)
1.3×10^{-6}	91.2	115.9
4.3×10^{-6}	101.5	109.9
7.5×10^{-6}	80.3	105.9
1.5×10^{-5}	120.2	117.1
6.3×10^{-5}	102.6	76.3
4×10^{-4}	87.2	67
3×10^{-3}	85	83.4
8×10^{-3}	68.7	90.9
1.5×10^{-2}	65.3	89.1

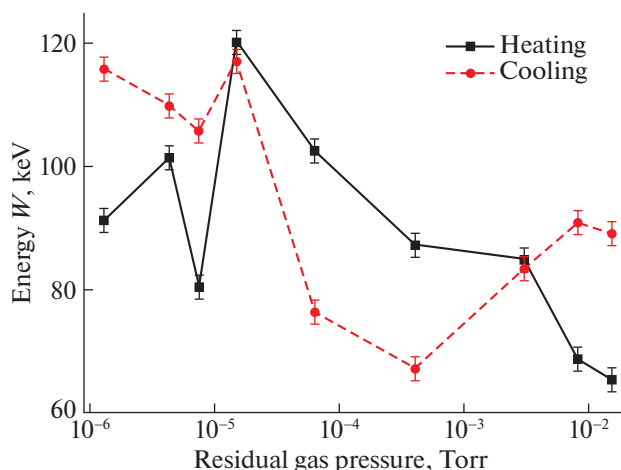
the pyroelectric deflector in vacuum at a residual gas pressure of 1.5×10^{-5} Torr. The X-ray spectrum measured upon cooling of the pyroelectric source at a residual gas pressure of 1.5×10^{-5} Torr is shown in Fig. 2b. Both spectra are presented on a log scale.

Both figures contain characteristic X-ray peaks with energies of 16.6, 18.6, and 6.4 keV corresponding to K_{α} and K_{β} lines of niobium whose atoms include into the pyroelectric crystal composition, and iron K_{α} , i.e., a vacuum chamber material. To the left of the iron K_{α} line, we can see a peak with an energy of ~ 1.5 keV corresponding to aluminum K_{α} line. The presence of this peak in the measured spectra confirms the feasibility of acceleration of electrons emitted from the pyroelectric crystal surface to grounded aluminum heat conductors. Over the main X-ray spectra, the same spectra on the linear scale are shown on the right top corner. Each additional spectrum in Fig. 2 shows the approximating straight line. The end-point energy of X-ray bremsstrahlung is defined by the abscissa of the intersection point of the approximating straight line to the horizontal axis.

Fitting was performed in the high-energy spectral range [5]. Using this method, the end-point energy of X-ray bremsstrahlung was determined in all spectra measured in this experiment. The results of processing of the obtained experimental data are listed in Table 1.

Figure 3 shows the dependences of the end-point energy (W , keV) of X-ray bremsstrahlung generated upon heating and cooling of the pyroelectric source at various residual gas pressures in the range from 1.3×10^{-6} Torr to 1.5×10^{-2} Torr.

The end-point energy of X-rays registered in experiments reaches a maximum of 120.2 keV in the case of heating (Fig. 2a) and 117.1 keV in the case of cooling (Fig. 2b) of the pyroelectric source at a residual gas pressure of 1.5×10^{-5} Torr. The absence of a unified shape of curves shown in Fig. 3 is presumably

**Fig. 3.** Dependences of the end-point energy of bremsstrahlung X-rays on the residual gas pressure.

associated with different temperature variation rate modes [3] during heating (the average value reaches $3^{\circ}\text{C}/\text{min}$) and cooling (the average value reaches $0.7^{\circ}\text{C}/\text{min}$) of the pyroelectric source.

Photons with energies to 150–175 keV are observed in the spectra presented in Fig. 2. The energy of these photons exceeds the end-point energy. The cause of the formation of such photons requires further investigation.

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CONFLICT OF INTEREST

The authors declare that they have no conflicts of interest.

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