# On the Index of the Dirichlet Problem for Elliptic Systems on the Plane 

A. P. Soldatov<br>Belgorod State University, Belgorod, Russia

## 1. STATEMENT OF THE PROBLEM

For the elliptic system

$$
\begin{equation*}
\frac{\partial^{2} u}{\partial y^{2}}-A_{1} \frac{\partial^{2} u}{\partial x \partial y}-A_{0} \frac{\partial^{2} u}{\partial x^{2}}=0 \tag{1}
\end{equation*}
$$

with constant matrix coefficients $A_{0}, A_{1} \in \mathbb{R}^{l \times l}$ in a bounded domain $D$, we consider the Dirichlet problem

$$
\begin{equation*}
\left.u\right|_{\partial D}=f \tag{2}
\end{equation*}
$$

This problem was analyzed by Bitsadze [1] for domains $D$ bounded by a smooth contour $\Gamma$. He also singled out the class of so-called weakly coupled systems (1); the Dirichlet problem for system (1) has the Fredholm property if and only if the system belongs to that class. In this case, the index of the problem proves to be zero. In the general case of second-order elliptic systems with continuous coefficients, these studies were continued in $[2-5]$ and other papers.

To system (1), one naturally assigns the quadratic matrix pencil

$$
P(z)=z^{2}-A_{1} z-A_{0}
$$

and the characteristic polynomial $\operatorname{det} P(z)$ of degree $2 l$. The ellipticity condition means that the roots of this polynomial do not lie on the real axis. Let $\sigma$ be the set of roots of this polynomial in the upper half-plane. Denoting the multiplicity of a root $\nu \in \sigma$ by $\lambda_{\nu}$, one can write

$$
\operatorname{det} P(z)=\prod_{\nu \in \sigma}(z-\nu)^{l_{\nu}}(z-\bar{\nu})^{l_{\nu}}, \quad \sum_{\nu \in \sigma} l_{\nu}=l .
$$

By [6], the weak coupling condition is equivalent to the invertibility of the matrix $\int_{-\infty}^{\infty} P^{-1}(t) d t$. Another equivalent definition of weak coupling was also given there. For any elliptic system (1), there exist two matrices $B, J \in \mathbb{C}^{l \times l}$ such that

$$
\operatorname{det}\left(\begin{array}{cc}
B & \bar{B}  \tag{3}\\
B J & \overline{B J}
\end{array}\right) \neq 0, \quad \operatorname{det}(z-J)=\prod_{\nu \in \sigma}(z-\nu)^{l_{\nu}}
$$

and the matrix relation

$$
\begin{equation*}
B J^{2}-A_{1} B J-A_{0} B=0 \tag{4}
\end{equation*}
$$

is valid. Here the matrix $J$ is defined modulo similarity and can also be chosen in Jordan form.
In the present paper, we assume that system (1) is weakly coupled and consider the Dirichlet problem in a domain $D$ bounded by a piecewise smooth contour $\Gamma$. Numerous studies (e.g., see [7-10], where a detailed bibliography is given) deal with elliptic problems in piecewise smooth
domains on the plane. The forthcoming considerations will be carried out in modified Hölder weighted spaces. Recall the definition [11] of these spaces.

On $\Gamma$, we choose a finite set $F$ containing all corners of the contour and consider the family $\lambda=\left(\lambda_{\tau}, \tau \in F\right)$ of real numbers. By $C_{\lambda}^{\mu}(\bar{D}, F), 0<\mu<1$, we denote the space of all continuous functions $\varphi(z)$ on $\bar{D} \backslash F$ such that

$$
\varphi_{*}(z)=\varphi(z) \prod_{\tau \in F}|z-\tau|^{\mu-\lambda_{\tau}} \in C^{\mu}(\bar{D}), \quad \varphi_{*}(\tau)=0, \quad \tau \in F
$$

This is a Banach space with respect to the transferred norm; it is a weighted space in the following sense: the space $C_{0}^{\mu}$ consists of bounded functions, and the multiplication by the weight function $\varrho(z)=\prod_{F}|z-\tau|^{\lambda_{\tau}}$ is an isomorphism of $C_{0}^{\mu}$ onto $C_{\lambda}^{\mu}$. By definition, if $\lambda=\mu$, then the space $C_{\mu}^{\mu}$ consists of functions $\varphi \in C^{\mu}(\bar{D})$ vanishing on $F$. The subspace $C_{\lambda}^{\mu}$ of continuously differentiable functions $\varphi$ on $\bar{D} \backslash F$ whose partial derivatives $\varphi_{x}$ and $\varphi_{y}$ belong to $C_{\lambda-1}^{\mu}$ will be denoted by $C_{\lambda}^{1, \mu}$.

Further, we introduce the extension of $C_{(\lambda)}^{\mu}(\bar{D}, F) \supseteq C_{\lambda}^{\mu}(\bar{D}, F)$ by polynomials $p(z)$ in the complex variable $z$. Obviously, this extension is finite-dimensional; consequently, $C_{(\lambda)}^{\mu}(\bar{D}, F)$ is a Banach space with respect to the corresponding norm. The space $C_{(\lambda)}^{1, \mu}$ has a similar meaning relative to $C_{\lambda}^{1, \mu}$. The functions $\varphi \in C_{(\lambda)}^{\mu}(\bar{D}, F)$ are continuous on $\bar{D} \backslash\left\{\tau \in F, \lambda_{\tau} \leq 0\right\}$ and behave as $O(1)|z-\tau|^{\lambda_{\tau}}$ as $z \rightarrow \tau$ if $\lambda_{\tau} \leq 0$ and as $\varphi(\tau)+O(1)|z-\tau|^{\lambda_{\tau}}$ if $\lambda_{\tau}>0$. Moreover, one can readily show that if $\lambda<1$, then any function $\varphi \in C^{1}(\bar{D} \backslash F)$ whose derivatives are $O(1)|z-\tau|^{\lambda_{\tau}-1}$ as $z \rightarrow \tau$ belongs to $C_{(\lambda)}^{\mu}(\bar{D}, F)$. Moreover, we have the inclusions

$$
C^{\mu}(\bar{D}) \subseteq C_{(\lambda)}^{\mu}(\bar{D}, F), \quad \lambda \leq \mu, \quad C^{1, \mu}(\bar{D}) \subseteq C_{(\lambda)}^{1, \mu}(\bar{D}, F), \quad \lambda \leq 1+\mu
$$

which become strict equalities for $\lambda=\mu$ and $\lambda=1+\mu$, respectively.
In a similar way, one can define all these spaces on a piecewise smooth contour $\Gamma$. Obviously, the restriction operator $\left.\varphi \rightarrow \varphi\right|_{\Gamma}$ is bounded: $C_{(\lambda)}^{\mu}(\bar{D}, F) \rightarrow C_{(\lambda)}^{\mu}(\Gamma, F)$, and similarly for the spaces $C_{(\lambda)}^{\mu}$. However, if we wish to have the same assertion for $C^{1, \mu}$, then the arcs forming the contour $\Gamma$ should be subjected to an additional Lyapunov condition. More precisely, each arc $\Gamma_{0} \subseteq \Gamma$ not containing points $\tau \in F$ as interior points should belong to the class $C^{1, \mu}$. As a result, for the pair $(\Gamma, F)$, we obtain the class denoted by $C^{1, \mu}$. For curves $\Gamma$ of this class, the restriction operator is bounded in the spaces $C_{\lambda}^{1, \mu}$ and $C_{(\lambda)}^{1, \mu}$.

Throughout the following, we consider the Dirichlet problem in the space $C_{(\lambda)}^{\mu}(\bar{D}, F)$ for $0<|\lambda|<1$ (i.e., for $0<\left|\lambda_{\tau}\right|<1$ for all $\tau \in F$ ) under the assumption that $\Gamma$ does not contain cusps and the pair ( $\Gamma, F)$ belongs to the class $C^{1, \mu+0}$ (i.e., $C^{1, \mu+\varepsilon}$ for some $\varepsilon>0$ ).

## 2. REDUCTION TO THE RIEMANN-HILBERT PROBLEM

Consider the first-order elliptic system

$$
\begin{equation*}
\frac{\partial \phi}{\partial y}-J \frac{\partial \phi}{\partial x}=0 \tag{5}
\end{equation*}
$$

with the matrix $J$ in (3), (4), whose solutions are referred to as $J$-analytic or hyperanalytic functions [12]. The main results of the theory of analytic functions remain valid for the solutions of this system; in this case, the matrix

$$
\begin{equation*}
(\alpha+i \beta)_{J}=\alpha 1+\beta J \tag{6}
\end{equation*}
$$

plays the role of the complex variable $\alpha+i \beta$, where 1 stands for the identity matrix. For example, the Taylor series expansion has the form

$$
\phi(z)=\sum_{k=0}^{\infty}\left(z-z_{0}\right)_{J}^{k} \phi^{(k)}\left(z_{0}\right), \quad \phi^{(k)}=\frac{\partial^{k} \phi}{\partial x^{k}}
$$

and the Cauchy integral is given by the relation

$$
\begin{equation*}
(I \varphi)(z)=\frac{1}{\pi i} \int_{\Gamma}(t-z)_{J}^{-1} d t_{J} \varphi(t) \tag{7}
\end{equation*}
$$

where the contour $\Gamma$ has a positive orientation with respect to $D$ and, by analogy with (6), we set

$$
(d \alpha+i d \beta)_{J}=1 d \alpha+J d \beta
$$

In this notation, any $J$-analytic function $\phi \in C(\bar{D})$ in $D$ can be represented by the Cauchy formula $2 \phi=I \varphi^{+}$. This formula remains valid for functions $\phi \in C_{\lambda}^{\mu}(\bar{D}, F),-1<\lambda<0$. In this case, the integral operator $I: C_{\lambda}^{\mu}(\bar{D}, F) \rightarrow C^{\mu}(\Gamma, F)$ is bounded, and we have the Sokhotskii-Plemelj formula

$$
\begin{equation*}
(I \varphi)^{+}\left(t_{0}\right)=\varphi\left(t_{0}\right)+(K \varphi)\left(t_{0}\right), \quad t_{0} \in \Gamma \backslash F \tag{8}
\end{equation*}
$$

Here $(K \varphi)\left(t_{0}\right)$ is a singular integral defined by analogy with $(7)$ for $z=t_{0}$.
For the elliptic system (1), the solutions of Eq. (5), where $J$ occurs in condition (3), play the same role as the analytic functions for the Laplace equation. This role is expressed by the representation

$$
\begin{equation*}
u=\operatorname{Re} B \phi \tag{9}
\end{equation*}
$$

of the general solution of system (1). The fact that the function (9) is a solution of system (1) readily follows from (4). Here the derivative $\phi^{\prime}$ is uniquely determined by $u$ modulo a constant term $\eta \in \mathbb{C}^{l}$. This formula can be inverted by differentiation with the use of the relations $2 u_{x}=B \phi^{\prime}+\overline{B \phi^{\prime}}$ and $2 u_{y}=B J \phi^{\prime}+\overline{B J \phi^{\prime}}$, which follow from (7). Then, in view of (3),

$$
\phi^{\prime}=2\left(C_{1} \frac{\partial u}{\partial x}+C_{2} \frac{\partial u}{\partial y}\right), \quad\left(\begin{array}{cc}
C_{1} & C_{2}  \tag{10}\\
\bar{C}_{1} & \bar{C}_{2}
\end{array}\right)=\left(\begin{array}{cc}
B & \bar{B} \\
B J & \overline{B J}
\end{array}\right)^{-1}
$$

In particular, $\phi^{\prime}$ is a single-valued function. However, the function $\phi$ itself, given by the curvilinear integral

$$
\phi(z)=\phi\left(z_{0}\right)+\int_{z_{0}}^{z} d t_{J}\left(C_{1} u_{x}+C_{2} u_{y}\right)(t)
$$

is not necessarily single-valued in a multiply connected domain. In general, the nature of its multivaluedness depends on the connectedness degree of the domain $D$, that is, on the number $s$ of connected components of the contour $\Gamma=\partial D$. In an implicit form, the corresponding result [12] can be stated as follows.

To be definite, we assume that $z=0 \in D$. Then there exist $l \times l$ matrices of functions $u_{j} \in C^{\infty}(\bar{D}), 1 \leq j \leq s$, whose columns satisfy Eq. (1) such that any solution of this equation can be uniquely represented in the form

$$
\begin{equation*}
u=\operatorname{Re} B \phi+\sum_{j=1}^{s} u_{j} \xi_{j}, \quad \phi(0)=0 \tag{11}
\end{equation*}
$$

with some $\xi_{j} \in \mathbb{R}^{l}$. The solution of the Dirichlet problem (1), (2) is constructed in the class of functions of this form, where $\phi \in C_{(\lambda)}^{\mu}(\bar{D}, F), 0<|\lambda|<1$.

By $m_{\lambda}$ we denote the number of points $\tau \in F$ such that $\lambda_{\tau}>0$; the case in which $\lambda<0$ is not excluded and corresponds to $m_{\lambda}=0$. One can readily see that any $J$-analytic function $\phi \in C_{(\lambda)}^{\mu}(\bar{D}, F)$ can be uniquely represented in the form

$$
\phi=\phi_{0}+\sum_{j=1}^{m_{\lambda}} z_{J}^{j} \eta_{j}, \quad \phi_{0}(0)=0
$$

with some $\eta_{j} \in \mathbb{C}^{l}$. The substitution of this expression into (11) provides a similar representation in the class $\phi \in C_{\lambda}^{\mu}$ with the only difference that $s$ should be replaced by $s+m_{\lambda}$. In turn, the substitution of the last representation into condition (2) reduces the Dirichlet problem to the equivalent boundary value problem

$$
\begin{equation*}
\operatorname{Re} B \phi^{+}+\sum_{j=1}^{s+m_{\lambda}} b_{j} \xi_{j}=f, \quad \phi(0)=0 \tag{12}
\end{equation*}
$$

for a $J$-analytic function $\phi \in C_{\lambda}^{\mu}(\bar{D}, F)$ and $\xi_{j} \in \mathbb{R}^{l}$, where we have set $b_{j}=\left.u_{j}\right|_{\Gamma}$. Since $u_{j} \in C^{1}(\bar{D})$ and the contour $\Gamma$ contains no cusps, it follows that the matrix functions $b_{j}$ satisfy the Lipschitz condition on $\Gamma$ and hence belong to the class $C_{(\lambda)}^{\mu}(\Gamma, F), \lambda<1$. It is also obvious that the multiplication operator $\varphi \rightarrow b_{j} \varphi$ is bounded in $C_{(\lambda)}^{\mu}$.

We choose a scalar function $\chi_{\tau} \in C^{\infty}$ that is identically equal to unity in a neighborhood of $\tau$ and vanishes in a neighborhood of $F \backslash \tau$. Then the operator

$$
P f=f-\sum_{\tau, \lambda_{\tau}>0} \chi_{\tau} f(\tau)
$$

is the projection of the space $C_{(\lambda)}^{\mu}$ onto $C_{\lambda}^{\mu}$, and problem (12) can be represented in the form of the system

$$
\begin{equation*}
\operatorname{Re} B \phi^{+}+\sum_{j=1}^{s+m_{\lambda}}\left(P b_{j}\right) \xi_{j}=P f, \quad \phi(0)=0, \quad \sum_{j=1}^{s+m_{\lambda}} b_{j}(\tau) \xi_{j}=f(\tau), \quad \lambda_{\tau}>0 \tag{13}
\end{equation*}
$$

## 3. MAIN RESULTS

The operator of system (13) acts in the spaces $C_{(\lambda)}^{\mu}(\bar{D}, F) \times\left(\mathbb{R}^{l}\right)^{s+m_{\lambda}} \rightarrow C_{(\lambda)}^{\mu}(\Gamma, F) \times\left(\mathbb{R}^{l}\right)^{m_{\lambda}}$; i.e., the system is a finite-dimensional extension of the Riemann-Hilbert problem

$$
\begin{equation*}
\operatorname{Re} B \phi^{+}=f^{0} \tag{0}
\end{equation*}
$$

in the class of $J$-analytic functions $\phi \in C_{\lambda}^{\mu}$. The latter problem was studied in detail in [12], and the results can be restated for (13), i.e., for the original Dirichlet problem (1), (2).

For a sufficiently small $r>0$, the domain $D_{\tau}=D \cap\{|z-\tau|<r\}$ is a curvilinear sector with vertex $\tau$, whose opening angle is denoted by $\theta_{\tau}$. It is assumed that the vertex $\tau$ is not a cusp of the contour $\Gamma$, which is expressed by the inequality $0<\theta_{\tau}<2 \pi$. The curve $\Gamma_{\tau}=\Gamma \cap\{|z-\tau|<r\}$ consists of two smooth arcs $\Gamma_{\tau \pm 0}$, which are the lateral sides of the sector $D_{\tau}$. The notation is introduced so as to ensure that the passage from $\Gamma_{\tau-0}$ to $\Gamma_{\tau+0}$ through the vertex $\tau$ is performed in the positive sense of the contour $\Gamma$, i.e., in the sense for which the domain $D$ lies to the left. Let $q_{\tau \pm 0} \in \mathbb{C}$ be the unit tangent vector of the $\operatorname{arc} \Gamma_{\tau \pm 0}$ at the point $\tau$ directed from $\tau$ to the other endpoint of this arc, so that

$$
\begin{equation*}
q_{\tau}=q_{\tau-0} q_{\tau+0}^{-1}=e^{i \theta_{\tau}} \tag{14}
\end{equation*}
$$

In accordance with notation (6), consider the matrix

$$
\begin{equation*}
Q_{\tau}=q_{\tau-0, J} q_{\tau+0, J}^{-1} \tag{15}
\end{equation*}
$$

by (3), its eigenvalues are the complex numbers

$$
\begin{equation*}
q_{\tau(\nu)}=q_{\tau-0, \nu} q_{\tau+0, \nu}^{-1}=\varrho_{\tau(\nu)} e^{i \theta_{\tau(\nu)}}, \quad \nu \in \sigma \tag{16}
\end{equation*}
$$

of multiplicity $l_{\nu}$. Here, by analogy with (6), $(\alpha+i \beta)_{\nu}$ is understood as the complex number $\alpha+\nu \beta$.
Since the points $\nu \in \sigma$ lie in the upper half-plane, it follows that the affine transformation $z \rightarrow z_{\nu}$ is invariant in this half-plane and leaves the points of the real axis fixed. In particular, it does not

## SOLDATOV

change the orientation of the plane. Therefore, by analogy with (14), the quantity $\theta_{\tau(\nu)}$ is the opening angle of the curvilinear sector $D_{\tau(\nu)}$ onto which $D_{\tau}$ is mapped under this transformation. Note that the relations $0<\theta_{\tau(\nu)}<\pi, \theta_{\tau(\nu)}=\pi$, and $\pi<\theta_{\tau(\nu)}<2 \pi$ are equivalent to $0<\theta_{\tau}<\pi$, $\theta_{\tau}=\pi$, and $\pi<\theta_{\tau}<2 \pi$, respectively. One can readily justify this fact by varying $\nu$ along the segment joining $\nu$ with the imaginary unit $i$.

To the matrix $Q_{\tau}$, we assign the complex powers $Q_{\tau}^{\zeta}=e^{\zeta \ln Q_{\tau}}$ and $\bar{Q}_{\tau}^{\zeta}=e^{\zeta \ln Q_{\tau}}$, which are treated in the sense of functions of matrices. The bar on the right-hand side of the second relation stands for complex conjugation. By using these powers, we introduce the family of entire matrix functions

$$
\begin{equation*}
X_{\tau}(\zeta)=\zeta^{-1}\left[Q_{\tau}^{\zeta}\left(B^{-1} \bar{B}\right)-\left(B^{-1} \bar{B}\right) \bar{Q}_{\tau}^{\zeta}\right], \quad X_{\tau}^{0}(\zeta)=\zeta^{-1}\left(Q_{\tau}^{\zeta}-\bar{Q}_{\tau}^{\zeta}\right) \tag{17}
\end{equation*}
$$

By (14), the eigenvalues of the matrix $X_{\tau}^{0}(\zeta)$ are $\zeta^{-1}\left(q_{\tau(\nu)}^{\zeta}-\bar{q}_{\tau(\nu)}^{\zeta}\right)$; they have the multiplicities $l_{\nu}$, $\nu \in \sigma$. Therefore,

$$
\begin{equation*}
\operatorname{det} X_{\tau}^{0}(\zeta)=\zeta^{-l} \prod_{\nu \in \sigma}\left(q_{\tau(\nu)}^{\zeta}-\bar{q}_{\tau(\nu)}^{\zeta}\right)^{l_{\nu}}=(2 i)^{l} \prod_{\nu \in \sigma}\left[\frac{\sin \theta_{\tau(\nu)} \zeta}{\zeta}\right]^{l_{\nu}} \tag{18}
\end{equation*}
$$

In particular, $\operatorname{det} X_{\tau}^{0}(\zeta) \neq 0$ in the strip $|\operatorname{Re} \zeta| \leq \delta$ for sufficiently small $\delta>0$.
In each strip $\alpha_{1} \leq \operatorname{Re} \zeta \leq \alpha_{2}$, we have $X_{\tau}(\zeta)\left(X_{\tau}^{0}\right)^{-1}(\zeta) \rightarrow B^{-1} \bar{B}$ uniformly with respect to $\operatorname{Re} \zeta$ as $\operatorname{Im} \zeta \rightarrow \infty$. Therefore, in this strip, the determinant det $X_{\tau}(\zeta)$ has finitely many zeros. The number of these zeros counting multiplicities in the open strip lying between the lines $\operatorname{Re} \zeta=0$ and $\operatorname{Re} \zeta=\alpha, \alpha \neq 0$, is denoted by $n_{\tau}(\alpha)$. The number $n_{\tau}^{0}$ has a similar meaning with respect to the line $\operatorname{Re} \zeta=0$. Next, we take a number $\delta>0$ small enough to ensure that $n_{\tau}(\delta)=n_{\tau}(-\delta)=0$ for all $\tau$ and introduce the integers

$$
\begin{equation*}
\Delta_{\tau}^{ \pm}=\left.\frac{1}{2 \pi} \arg \left(\frac{\operatorname{det} X_{\tau}}{\operatorname{det} X_{\tau}^{0}}\right)\right|_{ \pm \delta-i \infty} ^{ \pm \delta+i \infty} \tag{19}
\end{equation*}
$$

Note that

$$
\begin{equation*}
\Delta_{\tau}^{+}-\Delta_{\tau}^{-}=n_{\tau}^{0} \tag{20}
\end{equation*}
$$

by the Rouché theorem.
Theorem 1. The Dirichlet problem (1), (2) has the Fredholm property in the class $C_{(\lambda)}^{\mu}$, $0<|\lambda|<1$, if and only if

$$
\begin{equation*}
\operatorname{det} X_{\tau}(\zeta) \neq 0, \quad \operatorname{Re} \zeta=\lambda_{\tau}, \quad \tau \in F \tag{21}
\end{equation*}
$$

and its index is given by the formula

$$
\begin{equation*}
\varkappa=-\sum_{\tau \in F}\left[\Delta_{\tau}^{ \pm} \pm n_{\tau}\left(\lambda_{\tau}\right)\right] \tag{22}
\end{equation*}
$$

where the sign $\pm$ coincides with the sign $\operatorname{sgn} \lambda_{\tau}$.
Proof. Obviously, the operators of problems (13) and (13 ${ }^{\circ}$ ) are Fredholm equivalent, and their indices $\varkappa$ and $\varkappa^{0}$ are related by the formula

$$
\begin{equation*}
\varkappa=\varkappa^{0}+(s-2) l, \tag{23}
\end{equation*}
$$

where we have used the condition $\phi(0)=0$ in problem (13).
On the other hand, a criterion for the Fredholm property of problem ( $13^{0}$ ) and a formula for its index were given in [12]. Let $m$ be the number of elements of the set $F$. By [12], the end symbol of problem $\left(13^{0}\right)$ is the block diagonal $m \times m$ matrix with diagonal blocks

$$
X(\zeta ; \tau)=\left(\begin{array}{cc}
B & \bar{B}  \tag{24}\\
B Q_{\tau}^{\zeta} & \bar{B} \bar{Q}_{\tau}^{\zeta}
\end{array}\right)
$$

ON THE INDEX OF THE DIRICHLET PROBLEM FOR ELLIPTIC SYSTEMS
considered on the lines $\operatorname{Re} \zeta=\lambda_{\tau}$. The matrix $X^{0}(\zeta ; \tau)$ has the same meaning with $B$ replaced by the identity matrix. In this notation, the criterion for the Fredholm property of problem ( $13^{0}$ ) in the space $C_{\lambda}^{\mu}$ is given by the inequality

$$
\begin{equation*}
\operatorname{det} X(\zeta ; \tau) \neq 0, \quad \operatorname{Re} \zeta=\lambda_{\tau}, \quad \tau \in F \tag{25}
\end{equation*}
$$

and its index $\varkappa^{0}(\lambda)$ for $1 / 2 \leq \lambda<0$ is given by the formula

$$
\begin{equation*}
\varkappa^{0}(\lambda)=-\left.\sum_{\tau} \frac{1}{2 \pi i} \ln \frac{\operatorname{det} X(\zeta ; \tau)}{\operatorname{det} X^{0}(\zeta ; \tau)}\right|_{\lambda_{\tau}-i \infty} ^{\lambda_{\tau}+i \infty}+l(2-s) \tag{26}
\end{equation*}
$$

In the general case of an arbitrary weight order $\lambda$, this formula should be supplemented by the relation

$$
\begin{equation*}
\varkappa^{0}\left(\lambda^{\prime}\right)-\varkappa^{0}\left(\lambda^{\prime \prime}\right)=\sum_{\tau} k_{\tau}, \tag{27}
\end{equation*}
$$

where $k_{\tau}$ is the number of zeros of the function $\operatorname{det} X(\zeta ; \tau)$ in the strip between the lines $\operatorname{Re} \zeta=\lambda_{\tau}^{\prime}$ and $\operatorname{Re} \zeta=\lambda_{\tau}^{\prime \prime}$ with the plus sign if $\lambda_{\tau}^{\prime} \leq \lambda_{\tau}^{\prime \prime}$ and with the minus sign otherwise.

It follows from the comparison of the matrices (17) and (24) that

$$
\left(\begin{array}{cc}
B^{-1} & 0  \tag{28}\\
Q_{\tau}^{\zeta} B^{-1} & -B^{-1}
\end{array}\right) X(\zeta ; \tau)=\left(\begin{array}{cc}
1 & B^{-1} \bar{B} \\
0 & \zeta X_{\tau}(\zeta)
\end{array}\right)
$$

By setting $B=1$ here, one can write out a similar relation for the matrices $X^{0}\left(\zeta ; \tau_{0}\right)$ and $X_{\tau}^{0}(\zeta)$. Therefore, conditions (21) and (25) are equivalent, and formula (26) with $\lambda=-\delta$ in notation (19) can be represented in the form

$$
\varkappa^{0}(-\delta)=-\sum_{\tau} \Delta_{\tau}^{-}+l(2-s) .
$$

By the same argument, relation (27) applied to the weight orders $\lambda^{\prime}=\lambda$ and $\lambda^{\prime \prime}=-\delta$ gives

$$
\varkappa^{0}(\lambda)-\varkappa^{0}(-\delta)=\sum_{\tau} k_{\tau}, \quad k_{\tau}=\left\{\begin{array}{ccc}
n_{\tau}\left(\lambda_{\tau}\right) & \text { if } & \lambda_{\tau}<0 \\
-n_{\tau}\left(\lambda_{\tau}\right)-n_{\tau}^{0} & \text { if } & \lambda_{\tau}>0
\end{array}\right.
$$

By combining the last two relations with (20) and (23), we obtain (22), which completes the proof of the theorem.

Note that the integers (19) and hence the index (22) of the Dirichlet problem can be evaluated in a more explicit form if the function $\zeta^{l} \operatorname{det} X_{\tau}(\zeta)$ is even. Then, by (18), the ratio $\operatorname{det} X / \operatorname{det} X^{0}$ is also even; therefore, the quantities (19) have opposite signs. This, together with (20), implies that

$$
\begin{equation*}
\Delta_{\tau}^{ \pm}= \pm \frac{1}{2} n_{\tau}^{0} \tag{29}
\end{equation*}
$$

It follows from the same considerations that $n_{\tau}(\alpha)=-n_{\tau}(-\alpha)$; therefore, by (22), the indices of the Dirichlet problem in the classes $C_{( \pm \lambda)}^{\mu}$ have opposite signs.

Just as in [12], we supplement Theorem 1 with the corresponding result on the smoothness and asymptotics of the solution of the Dirichlet problem.

Theorem 2. If, under the assumptions of Theorem 1 , the function $u \in C_{(\lambda)}^{\mu}(\bar{D}, F)$ is a solution of the Dirichlet problem with right-hand side $f \in C_{(\lambda)}^{1, \mu}(\Gamma, F)$, then $u \in C_{(\lambda)}^{1, \mu}(\bar{D}, F)$.

Proof. As was mentioned above, the matrix functions $b_{j}$ belong to the class $C_{(\lambda)}^{1, \mu}(\Gamma, F)$. Consequently, if $\phi \in C_{(\lambda)}^{\mu}(\bar{D}, F)$ in problem (13), then the right-hand side $f^{0}=P f-\sum\left(P b_{j}\right) \xi_{j}$ of

## SOLDATOV

problem $\left(13^{0}\right)$ belongs to the class $C_{(\lambda)}^{1, \mu}(\Gamma, F)$. Therefore, it remains to apply the corresponding result [12] on the smoothness of the solution to this problem.

The family of spaces $C_{(\lambda)}^{\mu}$ is monotone decreasing (by inclusion) with respect to each of the parameters $\mu$ and $\lambda_{\tau}, \tau \in F$. Thus one can introduce the classes

$$
C_{(\lambda+0)}^{\mu}=\bigcup_{\varepsilon>0} C_{(\lambda+\varepsilon)}^{\mu}, \quad C_{(\lambda-0)}^{\mu}=\bigcap_{\varepsilon>0} C_{(\lambda-\varepsilon)}^{\mu} .
$$

If $\lambda=0$, then we denote them by $C_{(+0)}^{\mu}$ and $C_{(-0)}^{\mu}$ for brevity. As an example, by analogy with the complex powers $Q_{\tau}^{\zeta}$ contained in (17), in the curvilinear sector $D_{\tau}$, we consider the matrix functions $\ln (z-\tau)_{J}$ and $(z-\tau)_{J}^{\zeta_{0}}$. One can readily see that $(z-\tau)_{J}^{\zeta_{0}}\left[\ln (z-\tau)_{J}\right]^{j} \in C_{(\alpha-0)}^{\mu}$, $\operatorname{Re} \zeta_{0}=\alpha, j=0,1, \ldots$

Let the point $\zeta_{0}$ be a zero of the function $\operatorname{det} X_{\tau}(\zeta)$; thus the matrix function $X_{\tau}^{-1}(\zeta)$ has a pole of some order $r\left(\zeta_{0}\right) \geq 1$ at this point. Note that $r\left(\zeta_{0}\right)$ does not exceed the multiplicity of the zero $\zeta_{0}$ of the function $\operatorname{det} X_{\tau}(\zeta)$. To describe the asymptotic behavior of the solution $u(z)$ in the sector $D_{\tau}$ as $z \rightarrow \tau$, to each number $\alpha \in \mathbb{R}$, we assign the $J$-analytic function $\phi_{\tau}(z ; \alpha) \in C_{(\alpha-0)}^{\mu}$ of the form

$$
\begin{equation*}
\phi_{\tau}(z ; \alpha)=\sum_{\operatorname{Re} \zeta_{0}=\alpha} \sum_{j=0}^{r\left(\zeta_{0}\right)-1}(z-\tau)_{J}^{\zeta_{0}}\left[\ln (z-\tau)_{J}\right]^{j} c_{j}\left(\zeta_{0}\right) \tag{30}
\end{equation*}
$$

with some coefficients $c_{j}\left(\zeta_{0}\right) \in \mathbb{C}^{l}$, where the summation is performed over all zeros $\zeta_{0}$ of the function $\operatorname{det} X_{\tau}(\zeta)$ on the line $\operatorname{Re} \zeta_{0}=\alpha$. Strictly speaking, we use this expression only for $\alpha \neq 0$. If $\alpha=0$, then to the points $\zeta_{0}$, we add the point $\zeta_{0}=0$ and suppose that $r(0)=0$ for $\operatorname{det} X_{\tau}(0) \neq 0$. In this case, we set

$$
\begin{equation*}
\phi_{\tau}(z ; 0)=\sum_{\operatorname{Re} \zeta_{0}=0, \zeta_{0} \neq 0} \sum_{j=0}^{r\left(\zeta_{0}\right)-1}(z-\tau)_{J}^{\zeta_{0}}\left[\ln (z-\tau)_{J}\right]^{j} c_{j}\left(\zeta_{0}\right)+\sum_{j=0}^{r(0)}\left[\ln (z-\tau)_{J}\right]^{j} c_{j}(0) . \tag{0}
\end{equation*}
$$

Note that $\phi_{\tau}^{\prime}(z) \in C_{\alpha-1}^{\mu}\left(D_{\tau}, \tau\right)$ in both cases.
In the following assertion, we no longer assume that condition (21) is satisfied.
Theorem 3. Let $u$ be a solution of the Dirichlet problem in the class $C_{(\lambda-0)}^{\mu}(\bar{D}, F)$ [more precisely, the function $\phi$ occurring in the representation (12) belongs to this class], and let the righthand side $f$ of this problem belong to the class $C_{\left(\lambda_{\tau}+0\right)}^{\mu}\left(\Gamma_{\tau}, \tau\right)$ on the lateral boundary $\Gamma_{\tau}$ of the sector $D_{\tau}$ for some $\tau \in F$. Then there exists a function $\phi_{\tau}\left(z ; \lambda_{\tau}\right)$ of the form (30) or $\left(30^{\circ}\right)$ depending on the value of $\lambda_{\tau}$ such that

$$
\begin{equation*}
u_{\tau}=u-\operatorname{Re} B \phi_{\tau} \in C_{\lambda_{\tau}+0}^{\mu}\left(\bar{D}_{\tau}, \tau\right) . \tag{31}
\end{equation*}
$$

If, in addition, $f \in C_{\left(\lambda_{\tau}+0\right)}^{1, \mu}\left(\Gamma_{\tau}, \tau\right)$, then $u_{\tau} \in C_{\left(\lambda_{\tau}+0\right)}^{1, \mu}\left(\bar{D}_{\tau}, \tau\right)$.
Proof. Just as in Theorem 2, we consider problem ( $13^{0}$ ) with right-hand side $f^{0} \in C_{\lambda_{-}+0}^{\mu}\left(\Gamma_{\tau}, \tau\right)$. For its solution $\phi$, one can use the corresponding result [12] on the asymptotics. It is stated similarly to the assertion to be justified with the only difference that $r\left(\zeta_{0}\right)$ in (30) should be replaced by the order of the pole $\tilde{r}\left(\zeta_{0}\right)$ of the function $X^{-1}(\zeta ; \tau)$ at the point $\zeta_{0}$. Here there is no difference between the cases $\alpha \neq 0$ and $\alpha=0$. It follows from (28) that the quantities $r\left(\zeta_{0}\right)$ and $\tilde{r}\left(\zeta_{0}\right)$ are related by the formula

$$
\tilde{r}\left(\zeta_{0}\right)=r\left(\zeta_{0}\right)+\left(1-\operatorname{sgn}\left|\zeta_{0}\right|\right) .
$$

Thus we obtain all assertion of the theorem.
Corollary. Let $n_{\tau}^{0}=0$, i.e., $\operatorname{det} X_{\tau}(\zeta) \neq 0$ for $\operatorname{Re} \zeta=0$. Then each solution $u \in C_{-0}^{\mu}$ of the Dirichlet problem whose right-hand side satisfies, the condition $f \in C_{(+0)}^{\mu}\left(\Gamma_{\tau}, \tau\right)$ also belongs to the class $C_{(+0)}^{\mu}\left(\bar{D}_{\tau}, \tau\right)$ in the sector $D_{\tau}$.

## 4. SELF-ADJOINT ELLIPTIC SYSTEMS

We apply the obtained results to the Lagrange self-adjoint elliptic system

$$
\begin{equation*}
\sum_{i, j=1}^{2} A_{i j} \frac{\partial^{2} u}{\partial x_{i} \partial x_{j}}=0, \quad A_{i j}^{\mathrm{T}}=A_{j i} \tag{32}
\end{equation*}
$$

where $x_{1}=x, x_{2}=y$, and the symbol $T$ stands for matrix transposition. The Dirichlet problem (2) for this system is considered in the classes $C_{-0}^{\mu}=C_{(-0)}^{\mu}$ and $C_{(+0)}^{\mu}$. By Theorems 1 and 2, the class $V^{ \pm} \subseteq C_{( \pm 0)}^{\mu}$ of solutions of the homogeneous problem is finite-dimensional and lies in $C_{( \pm 0)}^{1, \mu}$. In particular, if $v \in V^{ \pm}$, then the conormal derivative

$$
\begin{equation*}
\frac{\partial v}{\partial \nu}=\sum_{i, j=1}^{2} n_{i} A_{i j} \frac{\partial u}{\partial x_{j}} \tag{33}
\end{equation*}
$$

where $n_{1}$ and $n_{2}$ are the components of the outward normal, belongs to the class $C_{-1 \pm 0}^{\mu}(\Gamma, F)$. Therefore, its product by a function in $C_{(\mp 0)}^{\mu}$ belongs to the class $C_{-1+0}^{\mu}(\Gamma, F)$ and hence is integrable. Here and throughout the following, the product $u v$ of two $l$-vectors is understood as the inner product $u_{1} v_{1}+\cdots+u_{l} v_{l}$.

Theorem 4. Let system (32) be weakly coupled, and let relation (29) be valid for all $\tau \in F$. Then the conditions

$$
\int_{\Gamma} f \frac{\partial v}{\partial \nu} d s=0, \quad v \in V^{\mp}
$$

are necessary and sufficient for the solvability of the problem in the class $C_{( \pm 0)}^{\mu}$. Accordingly, the index $\varkappa^{ \pm}$of the problem in this class is given by the relation

$$
\begin{equation*}
\varkappa^{ \pm}=\operatorname{dim} V^{ \pm}-\operatorname{dim} V^{\mp}=\mp \frac{1}{2} \sum_{\tau} n_{\tau}^{0} . \tag{35}
\end{equation*}
$$

Proof. Since system (32) is self-adjoint, we have the identity

$$
\sum_{i, j=1}^{2} \frac{\partial}{\partial x_{i}}\left[\left(A_{i j} \frac{\partial u}{\partial x_{j}}\right) v-u\left(A_{i j} \frac{\partial v}{\partial x_{j}}\right)\right]=0
$$

for arbitrary solutions $u$ and $v$. Consequently, if $u, v \in C_{(\neq 0)}^{1, \mu}$, then the bracketed expression in this identity belongs to the class $C_{-1+0}^{\mu}$. Therefore, one can use the Green formula, which implies that

$$
\int_{\Gamma}\left[\left(\frac{\partial u}{\partial \nu}\right) v-u \frac{\partial u}{\partial \nu}\right] d s=0
$$

in notation (33). In particular, condition (34 ${ }^{ \pm}$) is necessary for the solvability of problem (2), (32) in the class $C_{( \pm 0)}^{1, \mu}$. By the density argument, this fact remains valid for the class $C_{( \pm 0)}^{\mu}$.

Further, let us show that the number of linearly independent conditions ( $34^{ \pm}$) coincides with the dimension $\operatorname{dim} V^{\mp}$. In other words, $v \rightarrow \partial v / \partial \nu$ is a linear one-to-one correspondence in the class $V=V^{-}$.

Indeed, let $\partial v / \partial \nu=0$ on $\Gamma$ for some function $v \in V$. Then, by (34) and the relation $\left.v\right|_{\Gamma}=0$, for the partial derivatives $v_{x}$ and $v_{y}$ on $\Gamma$, we have the homogeneous system of linear equations

$$
\begin{equation*}
\left(n_{1} A_{11}+n_{2} A_{21}\right) v_{x}+\left(n_{1} A_{12}+n_{2} A_{22}\right) v_{y}=0, \quad-n_{2} v_{x}+n_{1} v_{y}=0 \tag{36}
\end{equation*}
$$

For the coefficient matrix of this system, we have

$$
\left(\begin{array}{cc}
\sum_{i} n_{i} A_{i 1} & \sum_{i} n_{i} A_{i 2} \\
-n_{2} & n_{1}
\end{array}\right)\left(\begin{array}{cc}
n_{1} & -n_{2} \\
n_{2} & n_{1}
\end{array}\right)=\left(\begin{array}{cc}
Q_{1} & Q_{2} \\
0 & 1
\end{array}\right)
$$

where for brevity, we have set $Q_{1}=\sum_{i, j} n_{i} n_{j} A_{i j}$ and $Q_{2}=\sum_{i} n_{i}\left(n_{1} A_{i 2}-n_{2} A_{i 1}\right)$. This relation, together with the ellipticity condition $\operatorname{det} Q_{1} \neq 0$ for Eq. (32), implies that system (36) has only the zero solution, i.e., $v_{x}=v_{y}=0$ on $\Gamma$. For $v$, we write out the representation (9), where, by (10), the derivative $\phi^{\prime}$ vanishes on $\Gamma$. It follows from the Cauchy theorem that $\phi^{\prime} \equiv 0$ in the domain $D$, and hence $v \equiv 0$.

Thus relations $\left(34^{ \pm}\right)$define $\operatorname{dim} V^{\mp}$ linearly independent orthogonality conditions. Theorem 1, together with (29), implies that the index $\varkappa^{ \pm}=\operatorname{dim} V^{ \pm}-k^{ \pm}$of the problem in the class $C_{( \pm)}^{\mu}$ is determined by the right-hand side of (35). In addition, by virtue of the preceding considerations, for the number $k^{ \pm}$of linearly independent solvability conditions for this problem, we have the lower bound $k^{ \pm} \geq \operatorname{dim} V^{\mp}$. Therefore, $0=\varkappa^{+}-\varkappa^{-}=\left(\operatorname{dim} V^{+}-k^{-}\right)+\left(\operatorname{dim} V^{-}-k^{+}\right)$, which is possible only if $k^{ \pm}=\operatorname{dim} V^{\mp}$. This readily implies all assertions of the theorem.

The absolute value of the indices (35) coincides with the dimension of the space $V^{0}$ in the expansion $V^{-}=V^{0} \oplus V^{+}$. By Theorem 3, each of the functions in a basis of the space $V^{0}$ admits the asymptotics (31) at the points $\tau \in F$ with functions $\phi_{\tau}$ of the form ( $30^{\circ}$ ) nonzero at least for some $\tau$. Obviously, the doubled dimension of the space $V^{0}$ is equal to the sum $\sum_{\tau} n_{\tau}^{0}$.

As was mentioned above, relation (29) is necessarily valid in the case of an even function $\zeta^{l} \operatorname{det} X_{\tau}(\zeta)$. In this case, under assumption (25), the indices $\varkappa^{ \pm}$of the problem in the spaces $C_{( \pm \lambda)}^{\mu}$ are also opposite. If, as before, $V^{ \pm} \subseteq C_{( \pm \lambda)}^{\mu}$ stands for the classes of solutions of the homogeneous problem, then, by Theorems 2 and $3, V^{+} \subseteq C_{(\lambda+0)}^{1, \mu}$ and $V^{-} \subseteq C_{(\lambda-0)}^{1, \mu}$, which implies that the solvability conditions ( $34^{ \pm}$) are meaningful in this case as well. By using the same reasoning, one can show that Theorem 4 is valid also in the spaces $C_{( \pm \lambda)}^{\mu}$.

## 5. SYSTEMS OF TWO EQUATIONS

If $l=2$, then the determinant of the matrices $X_{\tau}$ can be evaluated in closed form. We set

$$
\left(\begin{array}{ll}
B_{11} & B_{12}  \tag{37}\\
B_{21} & B_{22}
\end{array}\right)^{*}=\left(\begin{array}{rr}
B_{22} & -B_{12} \\
-B_{21} & B_{11}
\end{array}\right)=(\operatorname{det} B) B^{-1}
$$

In this notation, we have

$$
\begin{align*}
B^{*} \bar{B} & =\left(\begin{array}{cc}
r & -2 i p_{2} \\
2 i p_{1} & \bar{r}
\end{array}\right)  \tag{38}\\
r & =B_{22} \bar{B}_{11}-B_{12} \bar{B}_{21}, \quad p_{j}=\operatorname{Im}\left(B_{1 j} \bar{B}_{2 j}\right), \quad j=1,2, \\
|\operatorname{det} B|^{2} & =\operatorname{det} B^{*} \operatorname{det} \bar{B}=|r|^{2}-4 p_{1} p_{2} .
\end{align*}
$$

By (37), in definition (15) of the matrix $X_{\tau}$, one can replace $B^{-1}$ by $B^{*}$, which is assumed below. Therefore,

$$
\begin{equation*}
\zeta X_{\tau}(\zeta)=Q_{\tau}^{\zeta}\left(B^{*} \bar{B}\right)-\left(B^{*} \bar{B}\right) \bar{Q}_{\tau}^{\zeta} \tag{39}
\end{equation*}
$$

As was mentioned in Section 1, the matrix $J$ occurring in relations (3) and (4) can be chosen in Jordan form. We have only the following three cases for it:

$$
J_{I}=\left(\begin{array}{cc}
\nu_{1} & 0  \tag{40}\\
0 & \nu_{2}
\end{array}\right), \quad \nu_{1} \neq \nu_{2} ; \quad J_{I I}=\left(\begin{array}{cc}
\nu & 1 \\
0 & \nu
\end{array}\right) ; \quad J_{I I I}=\left(\begin{array}{cc}
\nu & 0 \\
0 & \nu
\end{array}\right)
$$

In the last case $J=J_{I I I}$, the matrices $Q^{\zeta}$ and $\bar{Q}^{\zeta}$ are scalar, and relation (39) can be reduced to the form $X_{\tau}(\zeta)=\zeta^{-1}\left(Q_{\tau}^{\zeta}-\bar{Q}_{\tau}^{\zeta}\right)\left(B^{*} \bar{B}\right)=X_{\tau}^{0}(\zeta)\left(B^{*} \bar{B}\right)$. In particular, the determinants of
the matrices $X$ and $X^{0}$ differ by a constant factor, whence it follows that $\Delta_{\tau}^{ \pm}=0$. Note that, in this case, system (1) can be reduced to a single scalar equation. Therefore, below we restrict our considerations to the first two cases in (40).

First, let $J=J_{I}$. In this case, the matrix $Q_{\tau}$ occurring in (15) has the diagonal form with diagonal entries (16). For brevity, we set $q_{j}=q_{\tau\left(\nu_{j}\right)}, j=1,2$, and in a similar way, we write $q_{j}=\varrho_{j} e^{i \theta_{j}}$. Then $q_{j}^{\zeta}=\varrho_{j}^{\zeta} e^{i \theta_{j} \zeta}, \bar{q}_{j}^{\zeta}=\varrho_{j}^{\zeta} e^{-i \theta_{j} \zeta}, Q_{\tau}^{\zeta}=\operatorname{diag}\left(q_{1}^{\zeta}, q_{2}^{\zeta}\right)$, and $\bar{Q} \bar{\tau}=\operatorname{diag}\left(\bar{q}_{1}^{\zeta}, \bar{q}_{2}^{\zeta}\right)$. This, together with (38) and (39), implies that

$$
\begin{aligned}
\zeta X_{\tau}(\zeta) & =\left(\begin{array}{cc}
r\left(q_{1}^{\zeta}-\bar{q}_{1}^{\zeta}\right) & -2 i p_{2}\left(q_{1}^{\zeta}-\bar{q}_{2}^{\zeta}\right) \\
2 i p_{1}\left(q_{2}^{\zeta}-\bar{q}_{1}^{\zeta}\right) & \bar{r}\left(q_{2}^{\zeta}-\bar{q}_{2}^{\zeta}\right)
\end{array}\right) \\
\zeta^{2} \operatorname{det} X_{\tau}(\zeta) & =|r|^{2}\left(q_{1}^{\zeta}-\bar{q}_{1}^{\zeta}\right)\left(q_{2}^{\zeta}-\bar{q}_{2}^{\zeta}\right)-4 p_{1} p_{2}\left(q_{1}^{\zeta}-\bar{q}_{2}^{\zeta}\right)\left(q_{2}^{\zeta}-\bar{q}_{1}^{\zeta}\right) \\
& =|\operatorname{det} B|^{2}\left(q_{1}^{\zeta}-\bar{q}_{1}^{\zeta}\right)\left(q_{2}^{\zeta}-\bar{q}_{2}^{\zeta}\right)-4 p_{1} p_{2}\left(q_{1}^{\zeta} \bar{q}_{2}^{\zeta}+q_{2}^{\zeta} \bar{q}_{1}^{\zeta}-q_{1}^{\zeta} \bar{q}_{1}^{\zeta}-q_{2}^{\zeta} \bar{q}_{2}^{\zeta}\right)
\end{aligned}
$$

A similar expression can be written out for $X^{0}$ with $r=\operatorname{det} B=1$ and $p_{j}=0$. After obvious transformations, we have

$$
\begin{align*}
-\frac{\zeta^{2} \operatorname{det} X_{\tau}(\zeta)}{4|\operatorname{det} B|^{2} \varrho_{1}^{\zeta} \varrho_{2}^{\zeta}} & =\sin \theta_{1} \zeta \sin \theta_{2} \zeta-\beta\left[\left(\frac{\varrho_{1}^{\zeta}}{\varrho_{2}^{\zeta}}-\frac{\varrho_{2}^{\zeta}}{\varrho_{1}^{\zeta}}\right)^{2}+4 \sin ^{2}\left(\frac{\theta_{1}-\theta_{2}}{2}\right) \zeta\right]  \tag{I}\\
\beta & =\frac{p_{1} p_{2}}{|\operatorname{det} B|^{2}}=\frac{\operatorname{Im}\left(B_{11} \bar{B}_{21}\right) \operatorname{Im}\left(B_{12} \bar{B}_{22}\right)}{|\operatorname{det} B|^{2}}
\end{align*}
$$

Note that here the quantity $\beta=\beta(B)$ is an invariant of system (1), i.e., is independent of the choice of the matrix $B$ in relations (3) and (4) for $J=J_{I}$. Indeed, by [13], any other matrix $\tilde{B}$ of this kind is related to $B$ by the formula $\tilde{B}_{i j}=\lambda_{j} B_{i j}, 1 \leq i, j \leq 2$, with some $\lambda_{j} \neq 0$. This implies that $\tilde{p}_{1} \tilde{p}_{2}=\left|\lambda_{1}\right|^{2}\left|\lambda_{2}\right|^{2} p_{1} p_{2}$ and $|\operatorname{det} \tilde{B}|^{2}=\left|\lambda_{1}\right|^{2}\left|\lambda_{2}\right|^{2}|\operatorname{det} B|^{2}$, whence it follows that $\beta(\tilde{B})=\beta(B)$. The expression for the determinant of the block matrix occurring in (3) was given in [13]. In the above-introduced notation, we have

$$
\frac{1}{2} \operatorname{det}\left(\begin{array}{cc}
B & \bar{B} \\
B J & \bar{B} \bar{J}
\end{array}\right)=\left|\nu_{1}-\nu_{2}\right|^{2}|\operatorname{det} B|^{2}\left(\beta-\frac{\operatorname{Im} \nu_{1} \operatorname{Im} \nu_{2}}{\left|\nu_{1}-\nu_{2}\right|^{2}}\right)
$$

Therefore, condition (3) is equivalent to the condition

$$
\begin{equation*}
\beta \neq \frac{\operatorname{Im} \nu_{1} \operatorname{Im} \nu_{2}}{\left|\nu_{1}-\nu_{2}\right|^{2}} \tag{I}
\end{equation*}
$$

For a given number $\beta \in \mathbb{R}$ satisfying this condition, one can choose an invertible matrix $B$ related to $\beta$ by the corresponding formula (41 $)$. In turn, by [13], this matrix defines a unique elliptic system (1) for which relations (3) and (4) are valid. In closed form, the coefficients $A_{0}$ and $A_{1}$ of this system can be found from the matrix relation

$$
\left(\begin{array}{cc}
0 & 1 \\
A_{0} & A_{1}
\end{array}\right)=\left(\begin{array}{cc}
B J & \bar{B} \bar{J} \\
B J^{2} & \bar{B} \bar{J}^{2}
\end{array}\right)\left(\begin{array}{cc}
B & \bar{B} \\
B J & \bar{B} \bar{J}
\end{array}\right)^{-1}, \quad J=J_{I}
$$

Now consider the case $J=J_{I I}$ in (40). In this case, for the matrices $q_{\tau \pm 0, J}$ occurring in (15), we have the expression

$$
q_{J}=q_{\nu}\left(\begin{array}{cc}
1 & q_{\nu}^{-1} \operatorname{Im} q \\
0 & 1
\end{array}\right)
$$

## SOLDATOV

Consequently,

$$
Q_{\tau}=q_{\tau(\nu)}\left(\begin{array}{cc}
1 & \delta_{\tau} \\
0 & 1
\end{array}\right), \quad \delta_{\tau}=\frac{\operatorname{Im} q_{\tau-0}}{q_{\tau-0}}-\frac{\operatorname{Im} q_{\tau+0}}{q_{\tau+0}}
$$

In notation (14), the complex number $\delta_{\tau}$ can be represented in the form

$$
\begin{equation*}
\delta_{\tau}=\left(\sin \theta_{\tau}\right) /\left(q_{\tau-0, \nu}\right)\left(q_{\tau-0, \nu}\right) . \tag{43}
\end{equation*}
$$

Therefore,

$$
Q_{\tau}^{\zeta}=q_{\tau(\nu)}^{\zeta}\left(\begin{array}{cc}
1 & \delta_{\tau} \zeta \\
0 & 1
\end{array}\right), \quad \bar{Q}_{\tau}^{\zeta}=\bar{q}_{\tau(\nu)}^{\zeta}\left(\begin{array}{cc}
1 & \bar{\delta}_{\tau} \zeta \\
0 & 1
\end{array}\right)
$$

This, together with (38) and (39), implies that

$$
\zeta X(\zeta)=\left(\begin{array}{cc}
r\left(q^{\zeta}-\bar{q}^{\zeta}\right)+2 i p_{1} \delta \zeta q^{\zeta} & -2 i p_{2}\left(q^{\zeta}-\bar{q}^{\zeta}\right)+\zeta\left(\bar{r} \delta q^{\zeta}-r \bar{\delta} \bar{q} \zeta\right) \\
2 i p_{1}\left(q^{\zeta}-\bar{q}^{\zeta}\right) & \bar{r}\left(q^{\zeta}-\bar{q}^{\zeta}\right)-2 i p_{1} \bar{\delta} \zeta \bar{q}^{\zeta}
\end{array}\right),
$$

where we have denoted $q=q_{\tau(\nu)}$ and $\delta=\delta_{\tau}$ for brevity. A similar expression is valid for the matrix $X^{0}$ with $r=1$ and $p_{j}=0$ as well. After simple manipulations, we obtain

$$
\begin{equation*}
-\frac{\zeta^{2} \operatorname{det} X_{\tau}(\zeta)}{4|\operatorname{det} B|^{2} \varrho^{2 \zeta}}=\sin ^{2} \theta \zeta-\beta|\delta|^{2} \zeta^{2}, \quad \beta=\frac{p_{1}^{2}}{|\operatorname{det} B|^{2}} \tag{II}
\end{equation*}
$$

By [13], the two matrices $\tilde{B}$ and $B$ occurring in (3) and (4) are related by the formula $\tilde{B}_{i 1}=\lambda_{1} B_{i 1}$, $\tilde{B}_{i 2}=\lambda_{2} B_{i 1}+\lambda_{1} B_{i 2}$ with some scalars $\lambda_{1} \neq 0$ and $\lambda_{2}$. Therefore, just as in the case of $J_{I}$, one can readily show that the quantity $\beta$ is an invariant of system (1). It also follows from [13] that

$$
\frac{1}{2} \operatorname{det}\left(\begin{array}{cc}
B & \bar{B} \\
B J & \bar{B} \bar{J}
\end{array}\right)=|\operatorname{det} B|^{2}\left(\beta-\operatorname{Im}^{2} \nu\right)
$$

in the case under consideration; therefore, condition (3) can be reduced to the form

$$
\begin{equation*}
\beta \neq \operatorname{Im}^{2} \nu \tag{II}
\end{equation*}
$$

If a given number $\beta$ satisfies condition $\left(42_{I I}\right)$ and the invertible matrix $B$ is related to it by the corresponding formula ( $41_{I I}$ ), then, just as above, one can show that this matrix defines a unique elliptic system with properties (3) and (4), where $J=J_{I I}$.

The determinant of the matrix $X^{0}$ satisfies the similar relations (41) with $\operatorname{det} B=1$ and $\beta=0$. These expressions imply that $\operatorname{det} X(\zeta)$ is an even function, whence we obtain (29). In this connection, the problem of evaluating the number $n_{\tau}^{0}$ of zeros of the function det $X(\zeta)$ on the line $\operatorname{Re} \zeta=0$ becomes important. By (41), these zeros coincide with the zeros of the function

$$
\begin{equation*}
|\operatorname{det} B|^{-2} \frac{\operatorname{det} X_{\tau}}{\operatorname{det} X_{\tau}^{0}}=1-\beta h, \tag{44}
\end{equation*}
$$

where we have set

$$
h_{I}(\zeta)=\frac{4}{\left(\sin \theta_{1} \zeta\right)\left(\sin \theta_{2} \zeta\right)}\left[\sinh ^{2}\left(\ln \frac{\varrho_{1}}{\varrho_{2}}\right) \zeta+\sin ^{2}\left(\frac{\theta_{1}-\theta_{2}}{2}\right) \zeta\right], \quad h_{I I}(\zeta)=\frac{\left|\delta_{\tau}\right|^{2} \zeta^{2}}{\sin ^{2} \theta_{\tau(\nu)} \zeta},
$$

respectively, for the two cases. [Recall that $\varrho_{j}=\varrho_{\tau\left(\nu_{j}\right)}$ and $\theta_{j}=\theta_{\tau\left(\nu_{j}\right)}$.] On the imaginary axis, we have

$$
h_{I}(i t)=\frac{4}{\left(\sinh t \theta_{1}\right)\left(\sinh t \theta_{2}\right)}\left[\sinh ^{2}\left(\frac{\theta_{1}-\theta_{2}}{2}\right) t-\sin ^{2}\left(\ln \frac{\varrho_{1}}{\varrho_{2}}\right) t\right], \quad h_{I I}(i t)=\frac{\left|\delta_{\tau}\right|^{2} t^{2}}{\sinh ^{2} t \theta_{\tau(\nu)}},
$$

and, in particular,

$$
\begin{equation*}
h_{I}(0)=\frac{4}{\theta_{1} \theta_{2}}\left[\left(\frac{\theta_{1}-\theta_{2}}{2}\right)^{2}-\left(\ln \frac{\varrho_{1}}{\varrho_{2}}\right)^{2}\right], \quad h_{I I}(0)=\frac{\left|\delta_{\tau}\right|^{2}}{\theta_{\tau(\nu)}^{2}} . \tag{45}
\end{equation*}
$$

Since $h(i t) \rightarrow 0$ as $t \rightarrow \infty$, we find that the equation $1-\beta h(i t)=0$ with $\beta h(0) \geq 1$ has at least one root on the half-line $t \geq 0$. Therefore, $n_{\tau}^{0} \geq 2$ for $\beta h(0) \geq 1$. This fact can be refined.

Lemma 1. If $J=J_{I I}$, then

$$
n_{\tau}^{0}= \begin{cases}2 & \text { if } \quad \beta h(0) \geq 1  \tag{46}\\ 0 & \text { if } \quad \beta h(0)<1 .\end{cases}
$$

A similar assertion is valid for $J=J_{I}$ under the additional assumption

$$
\begin{equation*}
\varrho_{\tau\left(\nu_{1}\right)}=\varrho_{\tau\left(\nu_{2}\right)}, \quad \theta_{\tau\left(\nu_{1}\right)}=3 \theta_{\tau\left(\nu_{2}\right)} . \tag{47}
\end{equation*}
$$

Proof. The function $g(x)=x^{2} / \sinh ^{2} x$ is strictly monotone decreasing from 1 to 0 on the half-line $[0, \infty)$. Since $1-\beta h(i t)=1-\beta h(0) g(t)$, we have the first assertion of the lemma.

Let $J=J_{I}$, and let relation (46) be valid; i.e., $\varrho_{1}=\varrho_{2}$ and $\theta_{1}=3 \theta_{2}$. Then

$$
h(i t)=\frac{4 \sinh t \theta_{1}}{\sinh 3 t \theta_{1}}=\frac{4}{3+4 \sinh ^{2} t \theta_{1}} .
$$

The function $g(x)=1 /\left(1+3 \sinh ^{2} x\right)$ is strictly monotone decreasing from 1 to 0 on the halfline $[0, \infty)$. Since $1-\beta h(i t)=1-\beta h(0) g\left(t \theta_{1}\right)$, we arrive at the second assertion of the lemma.

Note that under assumption (46), one can also readily write out the distribution of zeros of the function $\operatorname{det} X_{\tau}(\zeta)$ on the entire plane and, in particular, obtain a closed-form expression for the integer-valued characteristic $n_{\tau}(\alpha)$. Indeed, in this case, the expression (44 $)$ becomes much simpler:

$$
\zeta^{2} X_{\tau}(\zeta)=4|\operatorname{det} B|^{2} \varrho_{1}^{2 \zeta} \varrho_{2}^{2 \zeta}\left(\sin ^{2} \theta_{1} \zeta\right)\left(4 \sin ^{2} \theta_{1} \zeta+4 \beta \sin \theta_{1} \zeta-3\right) .
$$

Then we encounter the problem as to when relations (47) are valid, i.e., when there exist unit vectors $q_{\tau \pm 0}$ satisfying these relations. This question can be answered only partly.

Lemma 2. For given $\nu_{1} \neq \nu_{2}$, there exist unit vectors $q_{\tau \pm 0} \in \mathbb{C}$ such that the first relation in (47) is valid.

Proof. We have $q_{\tau \pm 0}=q_{\tau \pm 0}^{1}+i q_{\tau \pm 0}^{2}, q_{\tau \pm 0}^{j} \in \mathbb{R}$. Then

$$
q_{k}=\frac{q_{\tau-0}^{1}+\nu_{k} q_{\tau-0}^{2}}{q_{\tau+0}^{1}+\nu_{k} q_{\tau+0}^{1}}, \quad k=1,2 .
$$

The problem is to choose $q_{\tau \pm 0}$ such that $\left|q_{1}\right|=\left|q_{2}\right|$. This condition is equivalent to the relation

$$
\frac{q_{\tau-0}^{1}+\nu_{1} q_{\tau-0}^{2}}{q_{\tau-0}^{1}+\nu_{2} q_{\tau-0}^{1}}=\frac{q_{\tau+0}^{1}+\nu_{1} q_{\tau+0}^{2}}{q_{\tau+0}^{1}+\nu_{2} q_{\tau+0}^{1}}
$$

or

$$
\begin{equation*}
\frac{1+\nu_{1} t_{-}}{1+\nu_{2} t_{-}}=\frac{1+\nu_{1} t_{+}}{1+\nu_{2} t_{+}}, \quad t_{ \pm}=\frac{q_{\tau}^{2} \pm 0}{q_{\tau}^{1} \pm 0} . \tag{48}
\end{equation*}
$$

On the plane $w$, we consider the image $L$ of the real axis $\mathbb{R}$ under the linear-fractional transformation

$$
\begin{equation*}
w=\frac{1+\nu_{1} t}{1+\nu_{2} t} . \tag{49}
\end{equation*}
$$

## SOLDATOV

Obviously, $L$ is a circle passing through the points $w=1$ and $w=\nu_{1} / \nu_{2}$. The point $w=0$ lies outside this circle, since the points $t=-1 / \nu_{j}, j=1,2$, lie in the upper half-plane, i.e., on one side of the axis $\mathbb{R}$. Without loss of generality, the numbering of $\nu_{j}$ can be chosen so as to ensure that $\operatorname{Im}\left(\nu_{1} / \nu_{2}\right)>0$. If

$$
R_{0}<R<R_{1} ; \quad R_{0}=\min _{L}|w|, \quad R_{1}=\max _{L}|w|
$$

then the circle $|w|=R$ meets $L$ at two points $w_{ \pm}$. If $t_{ \pm}$corresponds to $w_{ \pm}$under the transformation (49), then formula (48) defines unit vectors $q_{\tau \pm 0}$ with the desired property.

The Jordan matrix $J$ occurring in relations (3) and (4) is referred to as the Jordan form of system (1).

Theorem 5. For a given matrix $J=J_{I}$ or $J=J_{I I}$ occurring in (40), there exists a contour $\Gamma$ with the unique corner $\tau$ and an elliptic system (1) with Jordan form $J$ such that the index of the Dirichlet problem (1), (2) in the class $C_{(+0)}^{\mu}$ is negative.

Proof. Since the function $\operatorname{det} X_{\tau}(\zeta)$ is even, it follows from Theorem 1 that the index $\varkappa^{+}$of the Dirichlet problem in the class $C_{(+0)}^{\mu}$ is equal to $-n_{\tau}^{0} / 2$.

First, let $J=J_{I I}$. In this case, the quantity $h(0)$ occurring in (45) is positive, and one can choose a number $\beta>0$ from condition $\left(42_{I I}\right)$ and the condition $\beta h(0) \geq 1$. Let $B$ be an invertible matrix related to $\beta$ by formula $\left(41_{I I}\right)$. Then the corresponding elliptic system is the desired one.

If $J=J_{I}$, then, following Lemma 2 , we choose $q_{\tau \pm 0}$ and consider a contour $\Gamma$ for which these vectors are tangent at the point $\tau$ to the lateral sides of the sector $D_{\tau}$. Then, by (45), $h(0)$ is positive, and it remains to repeat the above-performed considerations.

Note that if $J=J_{I I}$, then the contour $\Gamma$ can be chosen arbitrarily. The negativeness of the index $\varkappa^{+}$in the theorem implies that, in the case of the uniqueness of the solution of the Dirichlet problem, its existence is possible only under additional orthogonality conditions imposed on the right-hand side $f$ of Eq. (2). This phenomenon takes place only in the case of a corner. If $\tau$ is a regular point, i.e., $q_{\tau-0}=-q_{\tau+0}$, then, in notation (41), we have $\varrho_{1}=\varrho_{2}, \theta_{1}-\theta_{2}=2 \pi$ for $J=J_{I}$, and $\delta_{\tau}=0$ for $J=J_{I I}$. In both cases, the second term on the right-hand side in relation (41) identically vanishes; consequently, $\varkappa^{+}=0$.

## 6. THE LAMÉ SYSTEM WITH A NEGATIVE PARAMETER

Consider system (32) with coefficients

$$
\begin{array}{ll}
A_{11}=\left(\begin{array}{cc}
\lambda+2 & 0 \\
0 & 1
\end{array}\right), & A_{12}=\left(\begin{array}{cc}
0 & \lambda \\
1 & 0
\end{array}\right) \\
A_{21}=\left(\begin{array}{cc}
0 & 1 \\
\lambda & 0
\end{array}\right), & A_{22}=\left(\begin{array}{cc}
1 & 0 \\
0 & \lambda+2
\end{array}\right) \tag{50}
\end{array}
$$

For a positive value of $\lambda$, it is the Lamé system of plane isotropic elasticity. However, it can also be considered for negative values of the parameter $\lambda$. The characteristic polynomial of this system is equal to $(\lambda+2)\left(1+z^{2}\right)^{2}$, and hence the ellipticity condition can be reduced to the inequality $\lambda \neq-2$. If $\lambda=-1$, then the system becomes the Laplace equation, and this case can be excluded from considerations. If $\lambda \neq-1$, then relations (3) and (4) are valid for the matrices

$$
B=\left(\begin{array}{rr}
1 & 0 \\
i & -\varkappa
\end{array}\right), \quad J=\left(\begin{array}{cc}
i & 1 \\
0 & i
\end{array}\right), \quad \varkappa=\frac{\lambda+3}{\lambda+1}
$$

In particular, if $\lambda=-3$, then the system in question is strongly coupled. Therefore, the conditions imposed on the parameter $\lambda$ can be reduced to the inequalities

$$
\begin{equation*}
\lambda \neq-1,-2,-3 \tag{51}
\end{equation*}
$$

In this case, conditions $\left(42_{I I}\right)$ and $\left(43_{I I}\right)$ can be combined into a single condition of the form $\beta=1 / \varkappa^{2} \neq 1$. Here the validity of the inequality is provided by the condition $\lambda \neq-2$. In the case under consideration, relations (43) and ( $45_{I I}$ ) acquire the form

$$
h(0)=\frac{\left|\delta_{\tau}\right|^{2}}{\theta_{\tau}^{2}}, \quad \delta_{\tau}=\frac{\sin \theta_{\tau}}{q_{\tau-0} q_{\tau+0}}
$$

where we have used the fact that if $\nu=i$, then $z \rightarrow z_{\nu}$ is the identity mapping. Therefore, $h(0)=\left(\sin \theta_{\tau}\right)^{2} / \theta_{\tau}^{2}$, and the expression (46) can be reduced to the form

$$
n_{\tau}^{0}=\left\{\begin{array}{lll}
2 & \text { if } & \left|\sin \theta_{\tau}\right| \geq|\varkappa| \theta_{\tau} \\
0 & \text { if } & \left|\sin \theta_{\tau}\right|<|\varkappa| \theta_{\tau}
\end{array}\right.
$$

The inequality $\left|\sin \theta_{\tau}\right| \geq|\varkappa| \theta_{\tau}$, that is, the inequality

$$
\begin{equation*}
\left|\frac{\lambda+3}{\lambda+1}\right| \leq \frac{\left|\sin \theta_{\tau}\right|}{\theta_{\tau}}, \quad 0<\theta_{\tau}<2 \pi, \quad \theta_{\tau} \neq \pi \tag{52}
\end{equation*}
$$

together with (51), singles out two disjoint intervals

$$
\begin{equation*}
-3 \leq \lambda<\lambda_{\tau}^{0}, \quad \lambda_{\tau}^{1} \leq \lambda<-2 \tag{53}
\end{equation*}
$$

of the parameter $\lambda$ in which $n_{\tau}^{0}=2$. For the remaining values of $\lambda$, we have $n_{\tau}^{0}=0$. By applying Theorem 4 to system (32), (50), we obtain the following assertion.

Theorem 6. Let $V^{ \pm}$be the class of all solutions $v \in C_{( \pm 0)}^{\mu}$ of the homogeneous Dirichlet problem for system (32), (50). Then conditions $\left(34^{ \pm}\right)$are necessary and sufficient for the solvability of the problem in the class $C_{( \pm 0)}^{\mu}$, and its index $\varkappa^{ \pm}$is given by the relation

$$
\varkappa^{ \pm}=\mp \sum_{\tau} \chi\left(\frac{\left|\sin \theta_{\tau}\right|}{\theta_{\tau}}-\left|\frac{\lambda+3}{\lambda+1}\right|\right), \quad \chi(t)= \begin{cases}1 & \text { if } t \geq 0 \\ 0 & \text { if } t<0\end{cases}
$$

Note that system (32), (50) is strongly elliptic for $\lambda>-2$. In this case, by (53), inequality (52) fails, and therefore, $\varkappa^{+}=\varkappa^{-}=0$. In other words, the negative index $\varkappa^{+}$can appear only if $\lambda<-2$. More precisely, if the contour $\Gamma$ has a unique corner $\tau$ and $\lambda_{\tau}^{0} \leq \lambda<-2$, then $\varkappa^{-}=\operatorname{dim} V^{-}-\operatorname{dim} V^{+}=1$. Consequently, there exists exactly one (to within proportionality) solution of the homogeneous Dirichlet problem, which admits the asymptotics (31) with a nonzero function $\phi_{\tau}$ of the form $\left(30^{\circ}\right)$ at the point $\tau$.

## REFERENCES

1. Bitsadze, A.V., Kraevye zadachi dlya ellipticheskikh uravnenia vtorogo poryadka (Boundary Value Problems for Second-Order Elliptic Equations), Moscow: Nauka, 1966.
2. Tovmasyan, N.E., Differ. Uraun., 1966, vol. 2, no. 1, pp. 3-23; no. 2, pp. 163-171.
3. Vol'pert, A.I., Tr. Mosk. Mat. Obs., 1961, vol. 10, pp. 41-87.
4. Sirazhudinov, M.M., Izv. Ross. Akad. Nauk. Ser. Mat., 1997, vol. 61, no. 5, pp. 137-176.
5. Burskii, V.P., Metody issledovaniya granichnykh zadach dlya obshchikh differentsial'nykh uravnenii (Methods for Boundary Value Problems for General Differential Equations), Kiev, 2002.
6. Soldatov, A.P., Differ. Uravn., 2003, vol. 39, no. 5, pp. 674-686.
7. Kondrat'ev, V.A., Tr. Mosk. Mat. Obs., 1967, vol. 16, pp. 202-292.
8. Kondrat'ev, V.A. and Oleinik, O.A., Uspekhi Mat. Nauk, 1983, vol. 38, no. 2, pp. 3-76.
9. Dauge, M. and Steux, J.L., J. Differential Equations, 1987, vol. 70, no. 1, pp. 93-113.
10. Nazarov, S.A. and Plamenevskii, B.A., Ellipticheskie zadachi v oblastyakh s kusochno-gladkoi granitsei (Elliptic Problems in Domains with Piecewise Smooth Boundary), Moscow, 1991.
11. Soldatov, A.P., Deffer. Uravn., 2001, vol. 37, no. 6, pp. 825-838.
12. Soldatov, A.P., Izv. Ross. Akad. Nauk, 1992, vol. 56, no. 3, pp. 566-604.
13. Soldatov, A.P., in Sovremennaya matematika i ee prilozheniya (Modern Mathematics and Its Applications), Moscow, 2004, vol. 15, pp. 142-199.
