On the Relationship between the Integrated Cosine Function and the Operator Bessel Function

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The weakening of conditions imposed on the solution operators of the Cauchy problem for abstract first- and second-order differential equations has led (see [1–3]) to the notion of integrated semigroup and integrated cosine function.

In the present paper, we derive formulas relating the integrated cosine function to the solution operator $Y_k(t)$ of the Cauchy problem

$$u''(t) + \frac{k}{t}u'(t) = Au(t), \qquad t > 0,$$
 (1)

$$u(0) = u_0, \qquad u'(0) = 0,$$
 (2)

for the Euler-Poisson-Darboux equation in a Banach space E. (Here k > 0 is a parameter.)

The operator function $Y_k(t)$ was introduced in [4] and named the operator Bessel function. The set of operators A for which problem (1), (2) is uniformly well posed will be denoted by G_k . Thus if $A \in G_k$, then problem (1), (2) has a unique solution, which continuously depends on the initial data; moreover, $u(t) = Y_k(t)u_0$, $u_0 \in D(A)$, and

$$||Y_k(t)|| \le Me^{\omega t}, \qquad M \ge 1, \qquad \omega \ge 0. \tag{3}$$

Note that the condition for problem (1), (2) to be uniformly well posed and the properties of the operator Bessel function $Y_k(t)$ were given in [4].

Next, recall the definition of integrated cosine function.

Definition 1. Let $\alpha > 0$. A one-parameter family $C_{\alpha}(t)$, $t \geq 0$, of bounded linear operators is called an α -times integrated cosine function if the following conditions are satisfied:

1.

$$2\Gamma(\alpha)C_{\alpha}(t)C_{\alpha}(s) = \int_{t}^{t+s} (t+s-r)^{\alpha-1}C_{\alpha}(r)dr - \int_{0}^{s} (t+s-r)^{\alpha-1}C_{\alpha}(r)dr + \int_{t-s}^{t} (r-t+s)^{\alpha-1}C_{\alpha}(r)dr + \int_{0}^{s} (r+t-s)^{\alpha-1}C_{\alpha}(r)dr, \qquad t > s > 0.$$

- 2. $C_{\alpha}(0) = 0$.
- 3. $C_{\alpha}(t)x$ is a continuous function of $t \geq 0$ for each $x \in E$.
- 4. There exist constants M > 0 and $\omega \ge 0$ such that

$$||C_{\alpha}(t)|| \le Me^{\omega t}, \qquad t \ge 0. \tag{4}$$

The generator A of an integrated cosine function $C_{\alpha}(t)$ is defined as follows: the domain D(A) is the set of elements $x \in E$ such that there exists an element $y \in E$ satisfying the relation

$$C_{\alpha}(t)x - \frac{t^{\alpha}}{\Gamma(\alpha+1)}x = \int_{0}^{t} (t-r)C_{\alpha}(r)y\,dr, \qquad t \ge 0,$$
(5)

where $\Gamma(\cdot)$ is the Euler gamma function; in this case, we set Ax = y.

Theorem 1. Let $\alpha > 1$, let an operator A be the generator of an α -times integrated cosine function $C_{\alpha}(t)$, and let $u_0 \in D(A)$. Then problem (1), (2) is uniformly well posed (i.e., $A \in G_k$), and the corresponding operator Bessel function can be represented in the form

$$Y_{2\alpha}(t)u_0 = \frac{2^{\alpha}\Gamma(\alpha+1/2)}{\sqrt{\pi}t^{\alpha}} \left(C_{\alpha}(t)u_0 - \frac{1}{t} \int_0^t P'_{\alpha-1}\left(\frac{s}{t}\right) C_{\alpha}(s)u_0 ds \right),\tag{6}$$

where $P_{\nu}(\cdot)$ is a spherical Legendre function.

Proof. Formula (6) can be obtained heuristically as follows. Consider an operator cosine function C(t) and set

$$IC(t) = \int\limits_0^t C(au) d au.$$

Then, by the formula [4] for a parameter shift in Eq. (1),

$$Y_k(t)u_0 = \frac{2t^{1-k}\Gamma(k/2+1/2)}{\sqrt{\pi}\Gamma(k/2)} \int_0^t (t^2 - s^2)^{m-1} \frac{d^m (I^m C(s)u_0)}{ds^m} ds$$
 (7)

for $k=2m, m \in \mathbb{N}$, or, after simple transformations,

$$Y_{k}(t)u_{0} = \frac{2(-1)^{m-1}t^{1-k}\Gamma(k/2+1/2)}{\sqrt{\pi}\Gamma(k/2)} \int_{0}^{t} \frac{d^{m-1}\left((t^{2}-s^{2})^{m-1}\right)}{ds^{m-1}} \frac{d\left(I^{m}C(s)u_{0}\right)}{ds} ds$$

$$= \frac{2^{m}\Gamma(m+1/2)}{\sqrt{\pi}t^{m}} \int_{0}^{t} P_{m-1}\left(\frac{s}{t}\right) \frac{d\left(I^{m}C(s)u_{0}\right)}{ds} ds$$

$$= \frac{2^{m}\Gamma(m+1/2)}{\sqrt{\pi}t^{m}} \left(I^{m}C(t)u_{0} - \frac{1}{t} \int_{0}^{t} P'_{m-1}\left(\frac{s}{t}\right)I^{m}C(s)u_{0} ds\right), \tag{8}$$

where $P_{m-1}(\cdot)$ is a Legendre polynomial.

In (8), we replace $m \in N$ by $\alpha > 1$, $I^mC(t)$ by $C_{\alpha}(t)$, and the Legendre polynomial $P_{m-1}(\cdot)$ by the spherical Legendre function $P_{\alpha-1}(\cdot)$. Let us show that the function $Y_{2\alpha}(t)u_0$ defined in (6) is a solution of problem (1), (2) for $k = 2\alpha$.

Let us compute the first and second derivatives of $Y_{2\alpha}(t)u_0$:

$$\begin{split} Y_{2\alpha}'(t)u_0 &= \frac{2^{\alpha}\Gamma(\alpha+1/2)}{\sqrt{\pi}} \left(-\frac{\alpha+P_{\alpha-1}'(1)}{t^{\alpha+1}} C_{\alpha}(t)u_0 + \frac{1}{t^{\alpha}} C_{\alpha}'(t)u_0 \right. \\ &\quad + \frac{1+\alpha}{t^{\alpha+2}} \int_0^t P_{\alpha-1}'\left(\frac{s}{t}\right) C_{\alpha}(s)u_0 ds + \frac{1}{t^{\alpha+3}} \int_0^t P_{\alpha-1}'\left(\frac{s}{t}\right) s C_{\alpha}(s)u_0 ds \right), \\ Y_{2\alpha}''(t)u_0 &= \frac{2^{\alpha}\Gamma(\alpha+1/2)}{\sqrt{\pi}} \left(\frac{P_{\alpha-1}''(1) + 2(\alpha+1)P_{\alpha-1}'(1) + \alpha^2 + \alpha}{t^{\alpha+2}} C_{\alpha}(t)u_0 \right. \\ &\quad - \frac{2P_{\alpha-1}'(1) + 2\alpha}{t^{\alpha+1}} C_{\alpha}'(t)u_0 + \frac{1}{t^{\alpha}} C_{\alpha}''(t)u_0 - \frac{(\alpha+1)(\alpha+2)}{t^{\alpha+3}} \int_0^t P_{\alpha-1}'\left(\frac{s}{t}\right) C_{\alpha}(s)u_0 ds \\ &\quad - \frac{2\alpha+4}{t^{\alpha+4}} \int_0^t P_{\alpha-1}''\left(\frac{s}{t}\right) s C_{\alpha}(s)u_0 ds - \frac{1}{t^{\alpha+5}} \int_0^t P_{\alpha-1}''\left(\frac{s}{t}\right) s^2 C_{\alpha}(s)u_0 ds \right). \end{split}$$

Then, after integration by parts, we obtain

$$Y_{2\alpha}''(t)u_{0} + \frac{2\alpha}{t}Y_{2\alpha}'(t)u_{0}$$

$$= \frac{2^{\alpha}\Gamma(\alpha+1/2)}{\sqrt{\pi}} \left(\frac{\alpha-\alpha^{2}}{t^{\alpha+2}}C_{\alpha}(t)u_{0} - \frac{P_{\alpha-1}'(1)}{t^{\alpha+1}}C_{\alpha}'(t)u_{0} + \frac{1}{\Gamma(\alpha-1)t^{2}}u_{0} + \frac{1}{t^{\alpha}}C_{\alpha}(t)Au_{0} \right)$$

$$+ \frac{\alpha^{2}-\alpha}{t^{\alpha+3}} \int_{0}^{t} P_{\alpha-1}'\left(\frac{s}{t}\right)C_{\alpha}(s)u_{0}ds + \frac{1}{t^{\alpha+2}} \int_{0}^{t} \left(\frac{s^{2}}{t^{2}}P_{\alpha-1}''\left(\frac{s}{t}\right) + \frac{2s}{t}P_{\alpha-1}'\left(\frac{s}{t}\right)\right)C_{\alpha}'(s)u_{0}ds \right).$$
(9)

By [5, p. 206], the spherical Legendre function $P_{\alpha-1}(s)$ is a solution of the equation

$$(1 - s^2) w''(s) - 2sw'(s) + \alpha(\alpha - 1)w(s) = 0;$$

consequently, the function $P_{\alpha-1}(s/t)$ satisfies the relation

$$\frac{s^2}{t^2}w''\left(\frac{s}{t}\right) + \frac{2s}{t}w'\left(\frac{s}{t}\right) = w''\left(\frac{s}{t}\right) + \alpha(\alpha - 1)w\left(\frac{s}{t}\right). \tag{10}$$

By taking into account (10) and by integrating by parts, from (9), we obtain

$$Y_{2\alpha}''(t)u_{0} + \frac{2\alpha}{t}Y_{2\alpha}'(t)u_{0} = \frac{2^{\alpha}\Gamma(\alpha+1/2)}{\sqrt{\pi}} \left(\frac{1}{\Gamma(\alpha-1)t^{2}}u_{0} + \frac{1}{t^{\alpha}}C_{\alpha}(t)Au_{0} - \frac{1}{t^{\alpha+1}} \int_{0}^{t} P_{\alpha-1}'\left(\frac{s}{t}\right) \left(\frac{s^{\alpha-2}}{\Gamma(\alpha-1)}u_{0} + C_{\alpha}(s)Au_{0}\right) ds \right)$$

$$= \frac{2^{\alpha}\Gamma(\alpha+1)}{\sqrt{\pi}} \left(\frac{1}{t^{\alpha}}C_{\alpha}(t)Au_{0} - \frac{1}{t^{\alpha+1}} \int_{0}^{t} P_{\alpha-1}'\left(\frac{s}{t}\right)C_{\alpha}(s)Au_{0} ds + \frac{u_{0}}{\Gamma(\alpha-1)t^{2}} - \frac{u_{0}}{\Gamma(\alpha-1)t^{\alpha+1}} \int_{0}^{t} s^{\alpha-2}P_{\alpha-1}'\left(\frac{s}{t}\right) ds \right).$$

$$(11)$$

We use [6, Eq. 1.12.1.15] to compute the integral

$$\begin{split} \int_{0}^{t} s^{\alpha - 2} P'_{\alpha - 1} \left(\frac{s}{t} \right) ds &= t^{\alpha - 1} \int_{0}^{t} \tau^{\alpha - 2} P'_{\alpha - 1}(\tau) d\tau \\ &= t^{\alpha - 1} \Biggl(\tau^{\alpha - 2} P_{\alpha - 1}(\tau) - (\alpha - 2) \int_{0}^{t} \tau^{\alpha - 3} P_{\alpha - 1}(\tau) d\tau \Biggr) \\ &= t^{\alpha - 1} \Biggl(\tau^{\alpha - 2} P_{\alpha - 1}(\tau) - \frac{\tau^{\alpha - 2}}{\alpha - 1} \left(\alpha \tau P_{\alpha}(\tau) - \left((2\alpha - 1)\tau^{2} - \alpha + 1 \right) P_{\alpha - 1}(\tau) \right) \Biggr) \\ &= \frac{t^{\alpha - 1}}{\alpha - 1} \left((2\alpha - 1)\tau^{\alpha} P_{\alpha - 1}(\tau) - \alpha \tau^{\alpha - 1} P_{\alpha}(\tau) \right); \end{split}$$

hence

$$\int_{0}^{t} s^{\alpha-2} P'_{\alpha-1}\left(\frac{s}{t}\right) ds = t^{\alpha-1}.$$
(12)

It follows from (11), (12), and (6) that

$$Y_{2lpha}^{\prime\prime}(t)u_0+rac{2lpha}{t}Y_{2lpha}^{\prime}(t)u_0=AY_{2lpha}(t)u_0.$$

Therefore, the function $Y_{2\alpha}(t)u_0$ is a solution of Eq. (1).

Note that the representation (6) (after the change of variables $s = t\tau$ in the integral), together with (4), implies the estimate

$$||Y_{2\alpha}(t)|| \le M_1 e^{\omega t}. \tag{13}$$

To show that the function $Y_{2\alpha}(t)u_0$ satisfies the initial condition (2), we use relation (5) and the integral 2.17.1.4 in [6] and rewrite the expression (6) in the form

$$Y_{2\alpha}(t)u_{0} = \frac{2^{\alpha}\Gamma(\alpha+1/2)}{\sqrt{\pi}t^{\alpha}} \int_{0}^{t} P_{\alpha-1}\left(\frac{s}{t}\right) C_{\alpha}'(s)u_{0}ds$$

$$= \frac{2^{\alpha}\Gamma(\alpha+1/2)}{\sqrt{\pi}t^{\alpha}} \int_{0}^{t} P_{\alpha-1}\left(\frac{s}{t}\right) \left(\frac{s^{\alpha-1}u_{0}}{\Gamma(\alpha)} + \int_{0}^{s} C_{\alpha}(\varrho)Au_{0}d\varrho\right) ds$$

$$= \frac{2^{\alpha}\Gamma(\alpha+1/2)}{\sqrt{\pi}t^{\alpha}} \left(\frac{1}{\Gamma(\alpha)} \int_{0}^{1} \tau^{\alpha-1}P_{\alpha-1}(\tau)u_{0}d\tau + \int_{0}^{t} P_{\alpha-1}\left(\frac{s}{t}\right) ds \int_{0}^{s} C_{\alpha}(\varrho)Au_{0}d\varrho\right)$$

$$= u_{0} + \frac{2^{\alpha}\Gamma(\alpha+1/2)}{\sqrt{\pi}t^{\alpha-1}} \int_{0}^{1} P_{\alpha-1}(\tau)d\tau \int_{0}^{t\tau} C_{\alpha}(\varrho)Au_{0}d\varrho.$$

$$(14)$$

Now the desired assertion follows from (5), since the last term in (14) is of the order of t^2 as $t \to 0$. We prove the uniqueness of the solution of problem (1), (2) by contradiction. Let $u_1(t)$ and $u_2(t)$ be two solutions of problem (1), (2). Consider the function $w(t,s) = f(Y_{2\alpha}(s)(u_1(t) - u_2(t)))$ of two variables $t, s \ge 0$, where f belongs to the adjoint space E^* . Obviously, it satisfies the equation

$$\frac{\partial^2 w}{\partial t^2} + \frac{2\alpha}{t} \frac{\partial w}{\partial t} = \frac{\partial^2 w}{\partial s^2} + \frac{2\alpha}{s} \frac{\partial w}{\partial s}, \qquad t, s > 0,$$

and the conditions

$$w(0,s) = \frac{\partial w(0,s)}{\partial t} = \frac{\partial w(t,0)}{\partial s} = 0.$$

By using the change of variables $t_1 = (t+s)^2/4$, $s_1 = (t-s)^2/4$, one can reduce [7, Sec. 5, item 3] the last problem to the problem whose uniqueness in the class of twice continuously differentiable functions for $t, s \ge 0$ was proved in [7, Sec. 5, item 2]. Moreover, the desired uniqueness is also contained in Theorem 6.1 in [8], where even a more general equation was considered.

Consequently, $w(t, s) \equiv 0$, and since the functional $f \in E^*$ is arbitrary, for s = 0, we obtain the relation $u_1(t) \equiv u_2(t)$, and the proof of the uniqueness is complete.

Thus the operator function $Y_{2\alpha}(t)$ satisfies inequality (13), and the function $Y_{2\alpha}(t)u_0$ is the unique solution of problem (1), (2); consequently, problem (1), (2) is uniformly well posed. The proof of the theorem is complete.

Remark 1. Theorem 1 remains valid for $\alpha = 1$. In this case, the proof is much simpler, and $Y_2(t)u_0 = (1/t)C_1(t)u_0$, where $C_1(t)$ can naturally be referred to as the operator sine function.

Theorem 2. Let $A \in G_k$, k > 0, and let $Y_k(t)$ be the corresponding operator Bessel function. Then the operator A is the generator of an integrated cosine function $C_n(t)$, where n is the least positive integer such that $2n \ge k$.

Proof. First, note [4] that the operator Bessel function $Y_{2n}(t)$ can be expressed via the operator Bessel function $Y_k(t)$ with the use of the formula for a parameter shift:

$$Y_m(t) = rac{2}{\mathrm{B}(k/2+1/2,m/2-k/2)} \int\limits_0^1 \left(1-s^2
ight)^{(m-k)/2-1} s^k Y_k(ts) ds, \qquad m>k,$$

where $B(\cdot, \cdot)$ is the Euler beta function.

Let $u_0 \in D(A^n)$. By Theorem 3 in [9], the function

$$Y_0(t)u_0 = \frac{t}{(2n-1)!!} \left(\frac{1}{t}\frac{d}{dt}\right)^n \left(t^{2n-1}Y_{2n}(t)u_0\right)$$
(15)

is the unique solution of the equation

$$u''(t) = Au(t), \qquad t > 0, \tag{16}$$

with the initial conditions (2).

By using the relation [10, Eq. (1.13)]

$$\left(\frac{1}{t}\frac{d}{dt}\right)^n \left(t^{2n-1}Y_{2n}(t)u_0\right) = \sum_{j=0}^n \frac{2^{n-j}C_n^j\Gamma(n+1/2)}{\Gamma(j+1/2)} t^{2j-1} \left(\frac{1}{t}\frac{d}{dt}\right)^j Y_{2n}(t)u_0$$

and the expression [4]

$$Y_k'(t)u_0 = \frac{t}{k+1}Y_{k+2}(t)Au_0$$

for the derivative of the operator Bessel function, we rewrite formula (15) as

$$Y_0(t)u_0 = \sum_{j=0}^n \frac{2^{n-j}C_n^j\Gamma(n+1/2)}{\Gamma(j+1/2)} t^{2j} Y_{2n+2j}(t) A^j u_0.$$
 (17)

This, together with (3), implies the estimate

$$||Y_0(t)u_0|| \le M_1 e^{\omega_1 t} \sum_{j=0}^n ||A^j u_0||, \qquad \omega_1 > \omega.$$

Therefore, problem (16), (2) is exponentially uniformly n-well posed. It follows from Theorem 1.3 in [11] that the operator A is the generator of an integrated cosine function $C_n(t)$. The proof of the theorem is complete.

Remark 2. By virtue of the uniqueness of the solution, we have $C_n(t)u_0 = I^nY_0(t)u_0$ and $u_0 \in D(A^n)$, where $Y_0(t)$ is given by (15). After integration by parts, we obtain

$$C_{n}(t)u_{0} = I^{n}Y_{0}(t)u_{0} = \frac{1}{(n-1)!(2n-1)!!} \int_{0}^{t} (t-s)^{n-1}s \left(\frac{1}{s}\frac{d}{ds}\right)^{n} \left(s^{2n-1}Y_{2n}(s)u_{0}\right) ds$$

$$= \frac{(-1)^{n}}{(n-2)!(2n-1)!!} \left(s^{2n-1}\left(\frac{1}{s}\frac{d}{ds}\right)^{n-2} \left(\frac{(t-s)^{n-2}}{s}\right) Y_{2n}(s)u_{0}\right|_{0}^{t}$$

$$- \int_{0}^{t} s^{2n} \left(\frac{1}{s}\frac{d}{ds}\right)^{n-1} \left(\frac{(t-s)^{n-2}}{s}\right) Y_{2n}(s)u_{0}ds$$

for $n \geq 2$; moreover, the operator function $I^nY_0(t)$, originally defined on the dense [4] set $D(A^n)$, can be extended to the entire space E. For example,

$$egin{align} C_1(t) &= t Y_2(t), \qquad C_2(t) = rac{t^2}{3} Y_4(t) + rac{1}{3} \int \limits_0^t s Y_4(s) ds, \ C_3(t) &= rac{t^3}{15} Y_6(t) + rac{t}{5} \int \limits_0^t s Y_6(s) ds, \ \end{array}$$

and the inverse formulas read

$$egin{align} Y_2(t) &= rac{1}{t} C_1(t), \qquad Y_4(t) = rac{3}{t^2} C_2(t) - rac{3}{t^3} \int \limits_0^t C_2(s) ds, \ Y_6(t) &= rac{15}{t^3} C_3(t) - rac{45}{t^5} \int \limits_0^t s C_3(s) ds. \ \end{array}$$

In conclusion, we recall the definition of integrated semigroup and show how to use it so as to weaken the conditions imposed on the operator A occurring in the problem [12]

$$v'(t) + \frac{k}{t}v(t) = Av(t) + \frac{k}{t}g, \qquad t > 0,$$
 (18)

$$\lim_{t \to 0} \left(t^k v(t) \right) = v_0. \tag{19}$$

Definition 2. Let $\alpha > 0$. A one-parameter family of bounded linear operators $T_{\alpha}(t)$, $t \geq 0$, is called an α -times integrated semigroup if the following conditions are satisfied:

- 1. $\Gamma(\alpha)T_{\alpha}(t)T_{\alpha}(s) = \int_{s}^{t+s} (t+s-r)^{\alpha-1}T_{\alpha}(r)dr \int_{0}^{t} (t+s-r)^{\alpha-1}T_{\alpha}(r)dr, \ t,s \ge 0.$
- 2. $T_{\alpha}(0) = 0$.
- 3. $T_{\alpha}(t)x$ is a continuous function of $t \geq 0$ for each $x \in E$.
- 4. There exist constants M > 0 and $\omega \in R$ such that $||T_{\alpha}(t)|| \leq Me^{\omega t}, t \geq 0$.

The generator A of an integrated semigroup $T_{\alpha}(t)$ is defined as follows: the domain D(A) is the set of elements $x \in E$ such that there exists an element $y \in E$ satisfying the relation

$$T_{\alpha}(t)x - \frac{t^{\alpha}}{\Gamma(\alpha+1)}x = \int_{0}^{t} T_{\alpha}(s)y\,ds, \qquad t \ge 0;$$
 (20)

in this case, we set Ax = y.

Theorem 3. Let $k \in N$, let A be the generator of a k-times integrated semigroup $T_k(t)$, and let $g \in D(A)$ and $v_0 \in D(A^{k+1})$. Then the function

$$v(t) = t^{-k} \left(k! \, T_k(t) g + T_k^{(k)}(t) v_0 \right)$$

is the unique solution of problem (18), (19), and moreover,

$$||v(t)|| \le Me^{\omega t} (||g|| + ||A^k v_0||) + \sum_{j=0}^k \frac{t^{j-k}}{j!} ||A^j v_0||.$$
 (21)

Proof. Note that the representation of the solution v(t) of problem (18), (19) was obtained from Theorem 7 in [12], where it was assumed that A is the generator of a C_0 -semigroup T(t), by replacing the kth fractional integral of the semigroup T(t) by the integrated semigroup $T_k(t)$.

By taking into account the relation (e.g., see [3])

$$T_k^{(k)}(t)v_0 = T_k(t)A^k v_0 + \sum_{j=0}^{k-1} \frac{t^j}{j!} A^j v_0$$
 (22)

and definition (20) of the generator of the integrated semigroup $T_k(t)$, we compute v'(t). We obtain

$$v'(t) = -kt^{-k-1} \left(k! \, T_k(t) g + T_k^{(k)}(t) v_0 \right) + t^{-k} \left(k! \, A T_k(t) g + kt^{k-1} g + A T_k^{(k)}(t) v_0 \right);$$

therefore,

$$v'(t) + \frac{k}{t}v(t) = Av(t) + \frac{k}{t}g;$$

i.e., the function v(t) satisfies Eq. (18).

It also follows from (20) and (22) that the initial condition is satisfied, since

$$\lim_{t \to 0} (t^k v(t)) = k! \lim_{t \to 0} T_k(t) g + \lim_{t \to 0} T_k^{(k)}(t) v_0 = v_0.$$

The estimate (21) is a consequence of item 4 of Definition 2 and relations (20) and (22). Indeed,

$$||v(t)|| \le t^{-k} \left(k! M_1 t^k e^{\omega t} ||g|| + M_1 t^k e^{\omega t} ||A^k v_0|| + \sum_{j=0}^{k-1} \frac{t^j}{j!} ||A^j v_0|| \right)$$

$$= M e^{\omega t} \left(||g|| + ||A^k v_0|| \right) + \sum_{j=0}^{k} \frac{t^{j-k}}{j!} ||A^j v_0||.$$

Finally, to prove the uniqueness, we note that the change of variables

$$v(t) = t^{-k}w(t) + k! t^{-k}T_k(t)q$$

reduces problem (18), (19) to the problem

$$w'(t) = Aw(t), \qquad w(0) = v_0,$$

which, by virtue of Theorem 1.2 in [3], has a unique solution. The proof of the theorem is complete.

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